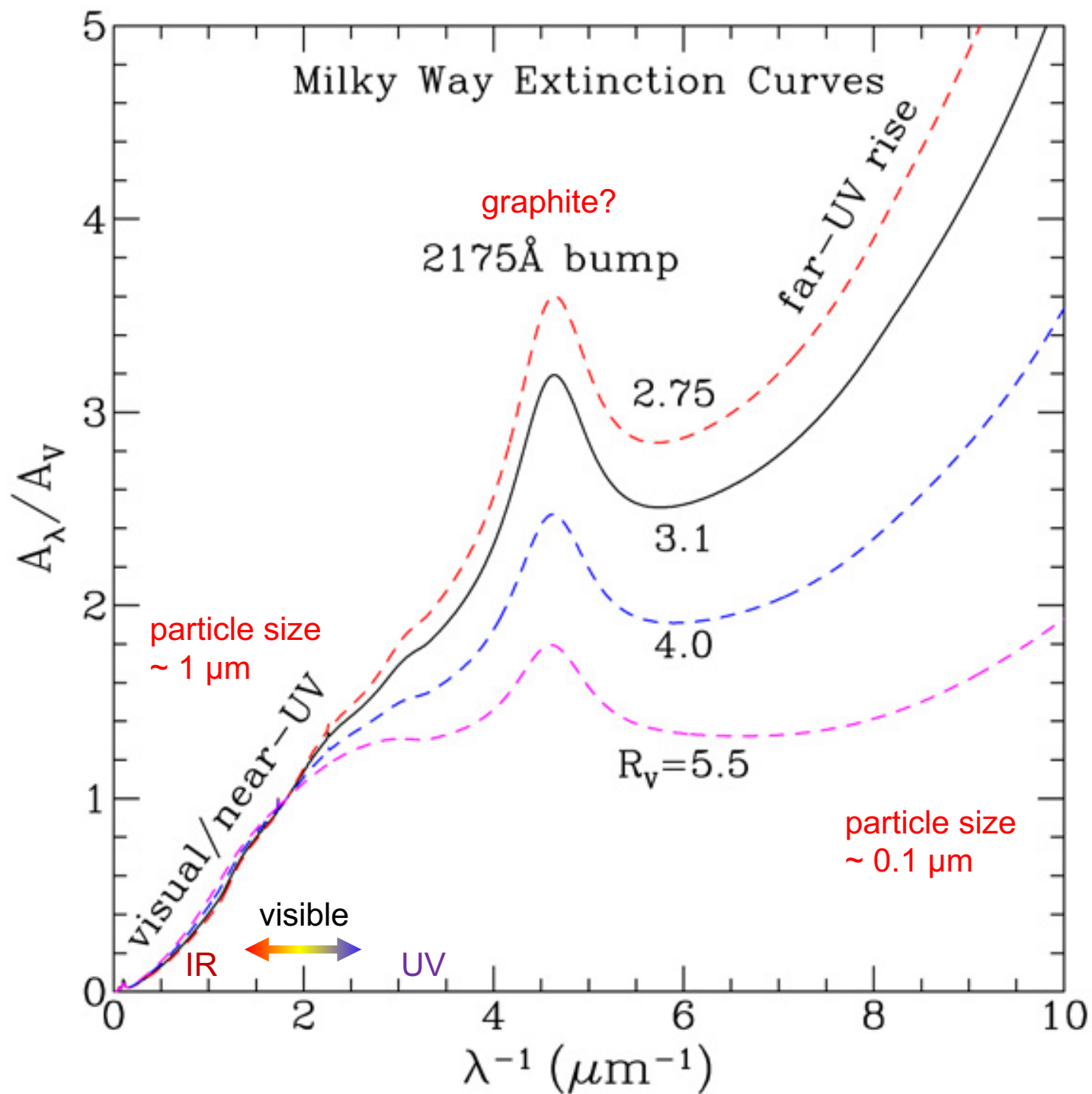
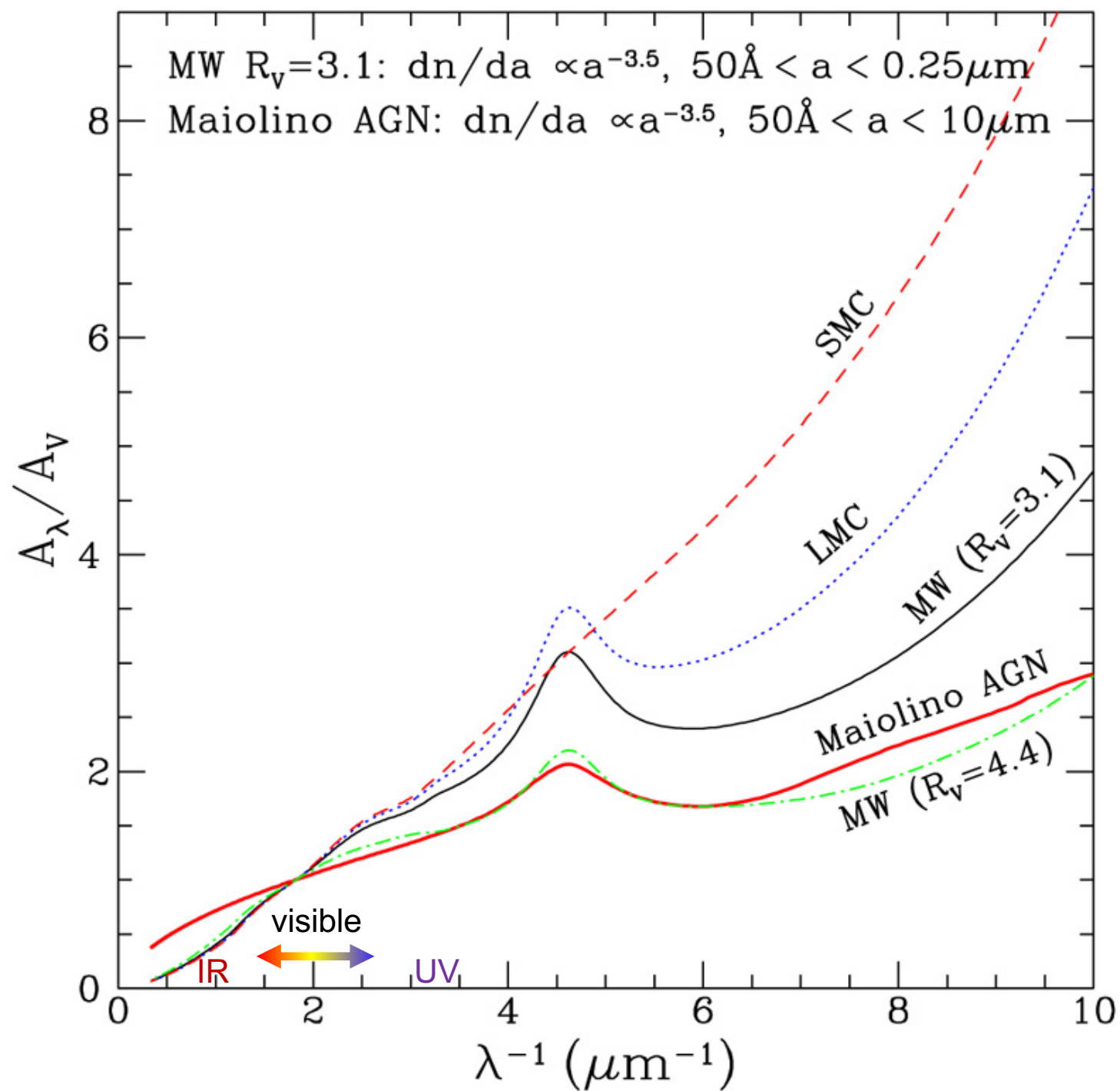


Interstellar Medium Phases

phase	component	T (K)	n (cm ⁻³)	M _{tot} (M _⊙)	scale (pc)
molecular cloud	cold	10 – 20	$> 10^3$	3×10^9	10 – 50
dust	cold/cool	—	$\sim 10^{-12} n_H$	few $\times 10^7$	
HI cloud	cool	100	20	6×10^9	100
HI cloud	warm	10^4	0.2		>1 kpc
corona/ bubble	hot	$> 10^6$	$\sim 10^{-3}$	10^8	galaxy



Li 2007 (AIP
Conference
Series)

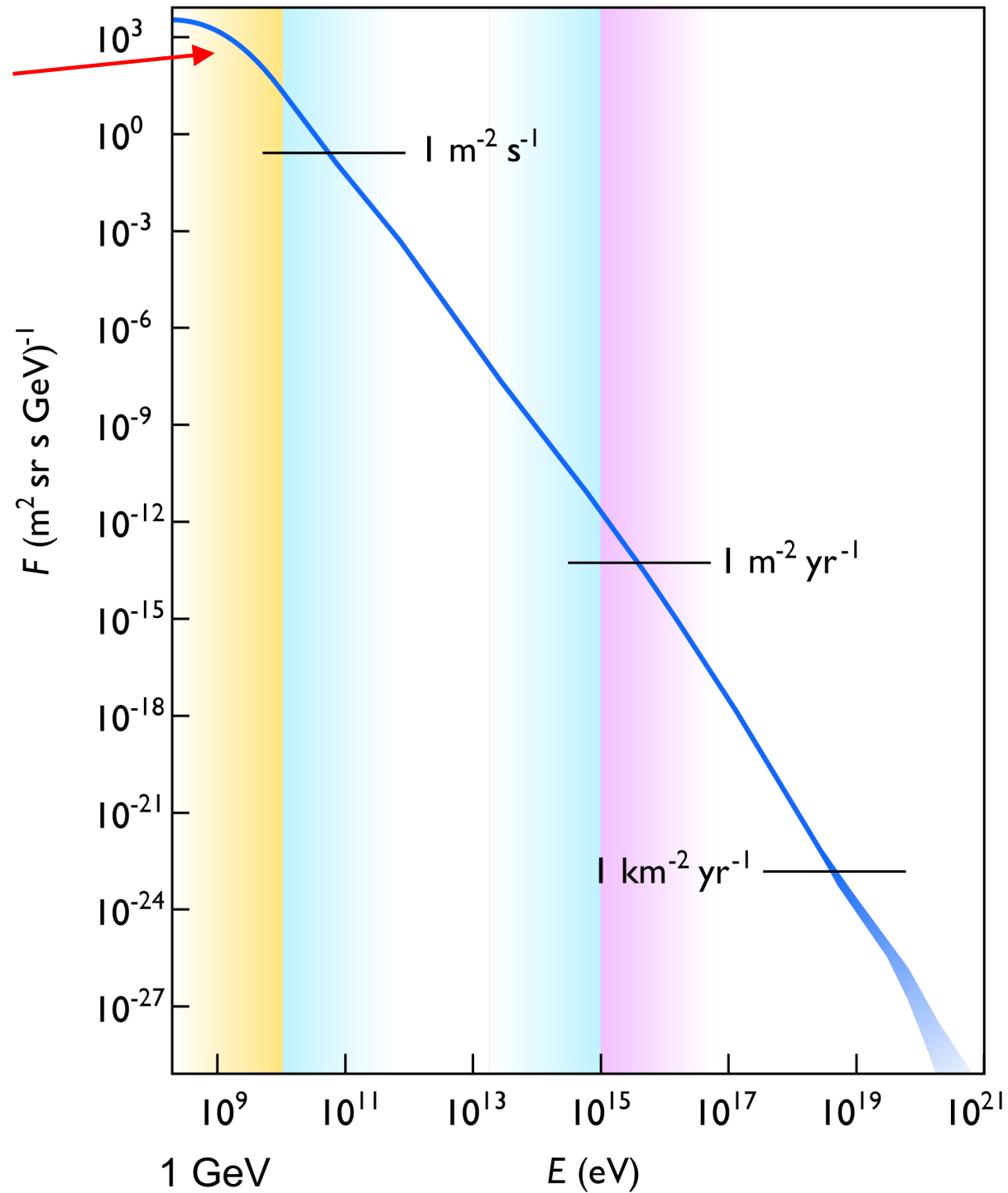


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Stellar Energy Budget

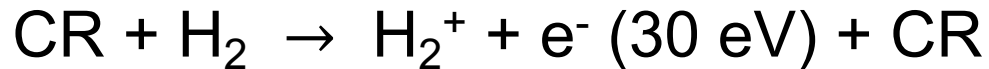
Source	Energy (J)
Stellar UV radiation, 20 $M_{\odot}$ star, lifetime ~ 3 Myr	4×10^{44}
Type II supernova	10^{44}
Wind, $\dot{M} \sim 5 \times 10^{-8} M_{\odot}/\text{yr}$, $v_{\infty} = 1000$ km/s, lifetime = 0.1 Myr	10^{46}
$10^3 M_{\odot}$ H cloud, $v \sim 10$ km/s, kinetic energy	10^{41}
$10^3 M_{\odot}$ H cloud, $T \sim 100$ K, thermal energy	2.5×10^{39}

interaction rate
 $\sim 2 \times 10^{-5} /s /H$

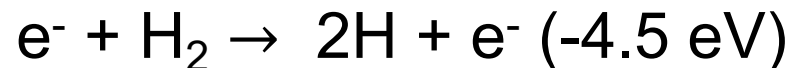


Sven Lafebre
2019
(Wikipedia)

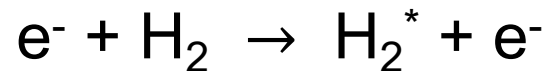
Cosmic ray heating



then



or



→ (Krumholz 2010) net yield over all possible interactions $\sim 7\text{--}20 \text{ eV}$

⇒ net heating rate $\Gamma_{\text{CR}} \sim 10^{-34} \text{ W per H nucleus}$

⇒ $\sim 10^{26} \text{ W}$ in a $1000 \text{ M}_\odot$ molecular cloud core

Cooling dense molecular gas

$T \sim 10\text{--}20\text{ K}$, $n > 10^3\text{ cm}^{-3}$

dust radiates as blackbody, $\lambda_{\text{max}} \sim 10\text{ }\mu\text{m}$ (optically thin)
but dust is only in effective thermal contact for $n > 10^4\text{ cm}^{-3}$

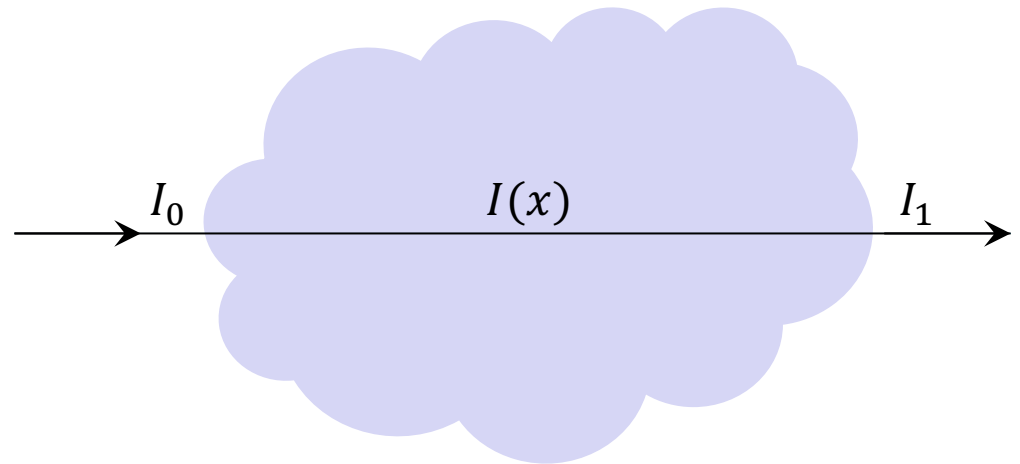
molecular lines are more effective at lower n — CO is most important

lines may be optically thick, limiting total emission

Radiative Transfer

transfer equation

$$\frac{dI}{dx} = j - \alpha I$$



often write $j = \rho\epsilon$, $\alpha = \rho\kappa$ (emissivity and opacity)

→ assume here constant j and α , $I = I_0$ for $x = 0$

(1) α negligible: solution is $I_1 = I_0 + \int j dx = I_0 + jx$

(2) j negligible: solution is $I_1 = I_0 e^{-\int \alpha dx} = I_0 e^{-\alpha x} = I_0 e^{-\tau}$

(3) general solution is $I_1 = I_0 e^{-\alpha x} + \frac{j}{\alpha} (1 - e^{-\alpha x})$
→ $\frac{j}{\alpha}$ as $\tau = \alpha x \rightarrow \infty$

effectively only see emission from within $\Delta x \sim 1/\alpha$ of surface

CO Cooling

Details depend on spatial and velocity structure; idealized calculation gives (Krumholtz 2010):

$$\Lambda_{CO} = x_{CO} (2J + 1) \frac{e^{-E_J/kT}}{Z} A_{J,J-1} (E_J - E_{J-1})$$

CO degeneracy Boltzmann decay line
fraction factor rate energy

where

$$E_J = hBJ(J + 1)$$

B = rotation constant

$$A_{J,J-1} = \frac{512\pi^4 B^3 \mu^2}{3hc^3} \frac{(J+1)^4}{2J+1}$$

μ = electric dipole moment

For $J = 5, T = 10 \text{ K}$, find

$$\Lambda_{CO} = 1.3 \times 10^{-34} \text{ W / H nucleus}$$

balance with CR heating for $T \sim 10 \text{ K}$, but quite temperature dependent

Table 2.5 Main processes that cool the interstellar gas

<i>Temperature</i>	<i>Cooling process</i>	<i>Spectral region</i>
$>10^7$ K	Free-free	X-ray
10^7 K $< T < 10^8$ K	Iron resonance lines	X-ray
10^5 K $< T < 10^7$ K	Metal resonance lines	UV, soft X-ray
8000 K $< T < 10^5$ K	C, N, O, Ne forbidden lines	IR, optical
Warm neutral gas: ~ 8000 K	Lyman- α , [OI]	1216 Å, 6300 Å
100 K $< T < 1000$ K	[OI], [CII], H ₂	Far IR: 63 μ m, 158 μ m
$T \sim 10-50$ K	CO rotational transitions	Millimeter-wave

Recombination

$$5000 \text{ K} < T < 20000 \text{ K}$$

$$\frac{dn_e}{dt} = -n_e^2 \alpha(T)$$

$$\alpha(T) = 2 \times 10^{-13} \left(\frac{T}{10^4 \text{ K}} \right)^{-3/4} \text{ cm}^3 \text{ s}^{-1}$$

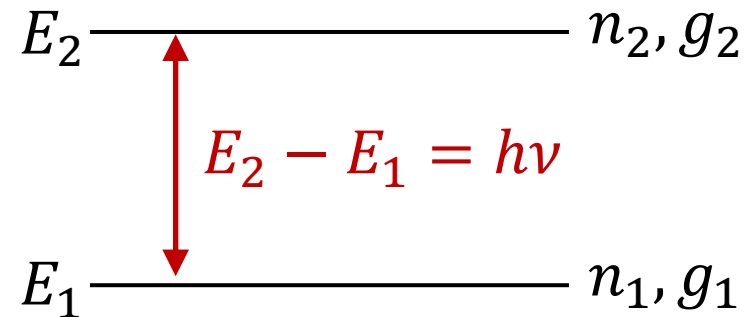
$$t_{rec} = \left| \frac{n_e}{dn_e/dt} \right| = 1500 \text{ yr} \left(\frac{T}{10^4 \text{ K}} \right)^{3/4} \left(\frac{n_e}{100 \text{ cm}^{-3}} \right)^{-1}$$

denser regions cool faster

Atomic Lines

$$\frac{dI_\nu}{dx} = j_\nu - \alpha_\nu I_\nu$$

2-level atom



$$\frac{dn_2}{dt} = -\underbrace{A_{21}n_2}_{\text{spontaneous emission}} + \underbrace{B_{12}u_\nu n_1}_{\text{absorption}} - \underbrace{B_{21}u_\nu n_2}_{\text{stimulated emission}}$$

$$\frac{A_{21}}{B_{21}} = \frac{8\pi h\nu^3}{c^3}$$

$$\frac{B_{21}}{B_{12}} = \frac{g_1}{g_2}$$

Atomic Lines

$$\frac{dI_\nu}{dx} = j_\nu - \alpha_\nu I_\nu$$

emission: $j_\nu = h\nu A_{21} n_2 \phi(\nu) / 4\pi$

absorption: $\alpha_\nu = \frac{c^2}{8\pi\nu^2} \frac{g_2}{g_1} n_1 A_{21} \phi(\nu) (1 - e^{-h\nu/kT})$

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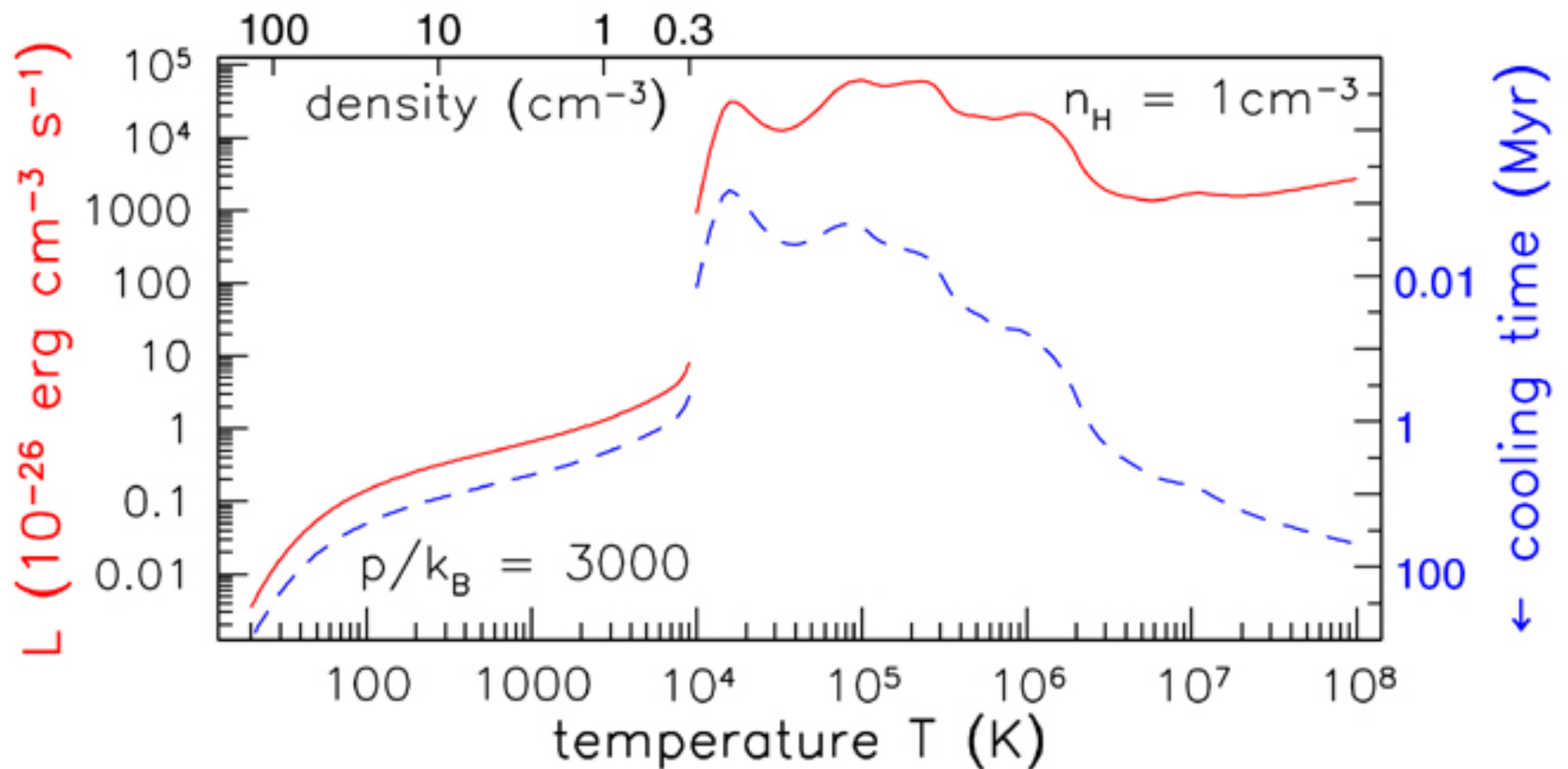


Fig 2.25 (Hensler, Wolfire) 'Galaxies in the Universe' Sparke/Gallagher CUP 2007