

- Eggleton, P.P. 1971, “The Evolution of Low Mass Stars,” *Monthly Notices of the Royal Astronomical Society*, vol. 151, pp. 351–364.
- London, R.A., and Flannery, B.P. 1982, “Hydrodynamics of X-Ray Induced Stellar Winds,” *Astrophysical Journal*, vol. 258, pp. 260–269.
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## 18.1 The Shooting Method

In this section we discuss “pure” shooting, where the integration proceeds from  $x_1$  to  $x_2$ , and we try to match boundary conditions at the end of the integration. In the next section, we describe shooting to an intermediate fitting point, where the solution to the equations and boundary conditions is found by launching “shots” from both sides of the interval and trying to match continuity conditions at some intermediate point.

Our implementation of the shooting method exactly implements multidimensional, globally convergent Newton-Raphson (§9.7). It seeks to zero  $n_2$  functions of  $n_2$  variables. The functions are obtained by integrating  $N$  differential equations from  $x_1$  to  $x_2$ . Let us see how this works.

At the starting point  $x_1$  there are  $N$  starting values  $y_i$  to be specified, but subject to  $n_1$  conditions. Therefore there are  $n_2 = N - n_1$  *freely specifiable* starting values. Let us imagine that these freely specifiable values are the components of a vector  $\mathbf{V}$  that lives in a vector space of dimension  $n_2$ . Then you, the user, knowing the functional form of the boundary conditions (18.0.2), can write a function or functor that generates a complete set of  $N$  starting values  $\mathbf{y}$ , satisfying the boundary conditions at  $x_1$ , from an arbitrary vector value of  $\mathbf{V}$  in which there are no restrictions on the  $n_2$  component values. In other words, (18.0.2) converts to a prescription

$$y_i(x_1) = y_i(x_1; V_0, \dots, V_{n_2-1}) \quad i = 0, \dots, N-1 \quad (18.1.1)$$

In the routine `Shoot` below, the function or functor that implements (18.1.1) will be called `load`, but you can pass it as an argument to the routine with any name of your choosing.

Notice that the components of  $\mathbf{V}$  might be exactly the values of certain “free” components of  $\mathbf{y}$ , with the other components of  $\mathbf{y}$  determined by the boundary conditions. Alternatively, the components of  $\mathbf{V}$  might parametrize the solutions that satisfy the starting boundary conditions in some other convenient way. Boundary conditions often impose algebraic relations among the  $y_i$ , rather than specific values for each of them. Using some auxiliary set of parameters often makes it easier to “solve” the boundary relations for a consistent set of  $y_i$ ’s. It makes no difference which way you go, as long as your vector space of  $\mathbf{V}$ ’s generates (through 18.1.1) all allowed starting vectors  $\mathbf{y}$ .

Given a particular  $\mathbf{V}$ , a particular  $\mathbf{y}(x_1)$  is thus generated. It can then be turned into a  $\mathbf{y}(x_2)$  by integrating the ODEs to  $x_2$  as an initial value problem (e.g., using Chapter 17’s `Odeint`). Now, at  $x_2$ , let us define a *discrepancy vector*  $\mathbf{F}$ , also of dimension  $n_2$ , whose components measure how far we are from satisfying the  $n_2$  boundary conditions at  $x_2$  (18.0.3). Simplest of all is just to use the right-hand sides

of (18.0.3),

$$F_k = B_{2k}(x_2, \mathbf{y}) \quad k = 0, \dots, n_2 - 1 \quad (18.1.2)$$

As in the case of  $\mathbf{V}$ , however, you can use any other convenient parametrization, as long as your space of  $\mathbf{F}$ 's spans the space of possible discrepancies from the desired boundary conditions, with all components of  $\mathbf{F}$  equal to zero if and only if the boundary conditions at  $x_2$  are satisfied. Below, you will be asked to supply a user-written function or functor that uses (18.0.3) to convert an  $N$ -vector of ending values  $\mathbf{y}(x_2)$  into an  $n_2$ -vector of discrepancies  $\mathbf{F}$ . Inside `Shoot`, this function is called `score`.

Now, as far as Newton-Raphson is concerned, we are nearly in business. We want to find a vector value of  $\mathbf{V}$  that zeros the vector value of  $\mathbf{F}$ . We do this by invoking the globally convergent Newton's method implemented in the routine `newt` of §9.7. Recall that the heart of Newton's method involves solving the set of  $n_2$  linear equations

$$\mathbf{J} \cdot \delta \mathbf{V} = -\mathbf{F} \quad (18.1.3)$$

and then adding the correction back,

$$\mathbf{V}^{\text{new}} = \mathbf{V}^{\text{old}} + \delta \mathbf{V} \quad (18.1.4)$$

In (18.1.3), the Jacobian matrix  $\mathbf{J}$  has components given by

$$J_{ij} = \frac{\partial F_i}{\partial V_j} \quad (18.1.5)$$

It is not feasible to compute these partial derivatives analytically. Rather, each requires a *separate* integration of the  $N$  ODEs, followed by the evaluation of

$$\frac{\partial F_i}{\partial V_j} \approx \frac{F_i(V_0, \dots, V_j + \Delta V_j, \dots) - F_i(V_0, \dots, V_j, \dots)}{\Delta V_j} \quad (18.1.6)$$

This is done automatically for you in the functor `NRfdjac` that comes with `newt`. The only input to `newt` that you have to provide is the routine `vecfunc` that calculates  $\mathbf{F}$  by integrating the ODEs. Here is the appropriate routine, a functor called `Shoot`, that is to be passed as the actual argument in `newt`:

```
shoot.h  template <class L, class R, class S>
          struct Shoot {
Function for use with newt to solve a two-point boundary value problem by shooting.
    Int nvar;                Number of coupled ODEs.
    Doub x1,x2;              Start and end points.
    L &load;                  Supplies initial values for ODEs from v[0..n2-1].
    R &d;                      Supplies derivative information to the ODE integrator.
    S &score;                 Returns the n2 functions that ought to be zero to satisfy
                                the boundary conditions at x2.
    Doub atol,rtol;
    Doub h1,hmin;
    VecDoub y;
    Shoot(Int nvarr, Doub xx1, Doub xx2, L &loadd, R &dd, S &scoree) :
        nvar(nvarr), x1(xx1), x2(xx2), load(loadd), d(dd),
        score(scoree), atol(1.0e-14), rtol(atol), hmin(0.0), y(nvar) {}
    Routine for use with newt to solve a two-point boundary value problem for nvar coupled
    ODEs by shooting from x1 to x2. Initial values for the nvar ODEs at x1 are generated
    from the n2 input coefficients v[0..n2-1], using the user-supplied routine load.
    VecDoub operator() (VecDoub_I &v) {
    This is the functor used by newt. It integrates the ODEs to x2 using an eighth-order Runge-
    Kutta method with absolute and relative tolerances atol and rtol, initial stepsize h1, and
```

```

    minimum stepsize hmin. At x2 it calls the user-supplied routine score and returns the
    n2 functions that ought to be zero. newt uses a globally convergent Newton's method to
    adjust the values of v until the returned functions are in fact zero.
    h1=(x2-x1)/100.0;
    y=load(x1,v);
    Output out;           No output generated by Odeint.
    Odeint<StepperDopr853<R> > integ(y,x1,x2,atol,rtol,h1,hmin,out,d);
    integ.integrate();
    return score(x2,y);
}
};

```

Note that Shoot is templated on the load, right-hand side for Odeint, and score routines. In practice, you will almost always want to write these as functors rather than functions. This makes communicating the various parameters in the problem easy — just pass them as parameters in the constructors.

For some problems the initial stepsize  $\Delta V$  might depend sensitively upon the initial conditions. It is straightforward to alter load to compute a suggested stepsize h1 as a member variable and feed it first to Shoot and hence to NRfdjac when the Shoot object is passed to newt.

A complete cycle of the shooting method thus requires  $n_2 + 1$  integrations of the  $N$  coupled ODEs: one integration to evaluate the current degree of mismatch, and  $n_2$  for the partial derivatives. Each new cycle requires a new round of  $n_2 + 1$  integrations. This illustrates the enormous extra effort involved in solving two-point boundary value problems compared with initial value problems.

If the differential equations are *linear*, then only one complete cycle is required, since (18.1.3) – (18.1.4) should take us right to the solution. A second round can be useful, however, in mopping up some (never all) of the roundoff error.

As given here, Shoot uses the high-efficiency eighth-order Runge-Kutta method of §17.2 to integrate the ODEs, but any of the other methods of Chapter 17 could just as well be used.

You, the user, must supply Shoot with: (i) a function or functor load(x1,v) that returns the n-vector y[0..n-1] (satisfying the starting boundary conditions, of course), given the freely specifiable variables of v[0..n2-1] at the initial point x1; (ii) a function or functor score(x2,y) that returns the discrepancy vector f[0..n2-1] of the ending boundary conditions, given the vector y[0..n-1] at the end-point x2; (iii) a starting vector v[0..n2-1]; (iv) a function or functor, called d in the routine, for the ODE integration; and other obvious parameters as described in the header comment above.

In §18.4 we give a sample program illustrating how to use Shoot.

#### CITED REFERENCES AND FURTHER READING:

- Acton, F.S. 1970, *Numerical Methods That Work*; 1990, corrected edition (Washington, DC: Mathematical Association of America).
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