

WELCOME TO

PHYS 201

Winter Quarter 09

Brigita Urbanc  
brigita@drexel.edu

[www.physics.drexel.edu/~wking/  
course/index.shtml](http://www.physics.drexel.edu/~wking/course/index.shtml)



download syllabus

○ Recitations & Homework:

Tim Jones tdj28@drexel.edu

Trevor King wking@drexel.edu

Labs:

○ Rachael Kratzer rmk55@...

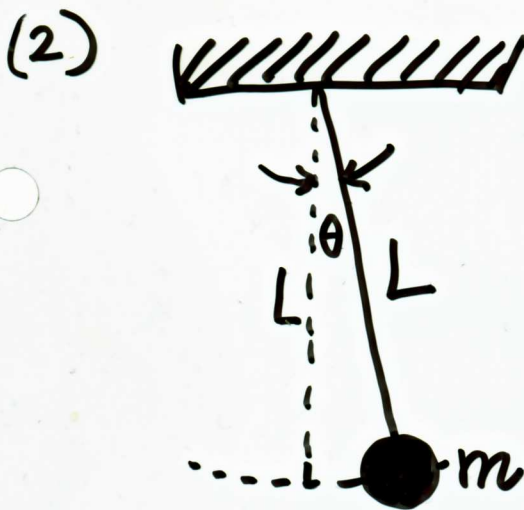
Pradeesh Nair prn25@...

# Oscillatory Motion



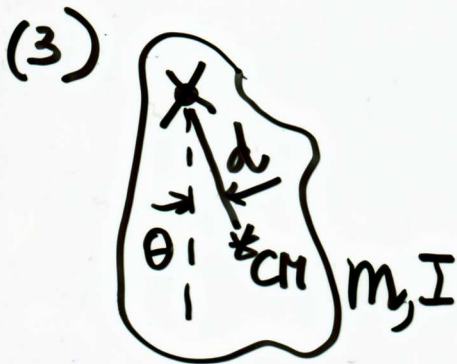
a particle attached to a string

↓  
simple harmonic oscillator



a simple pendulum

↓  
at small  $\theta \ll "1"$   
⇒ simple harmonic oscillator



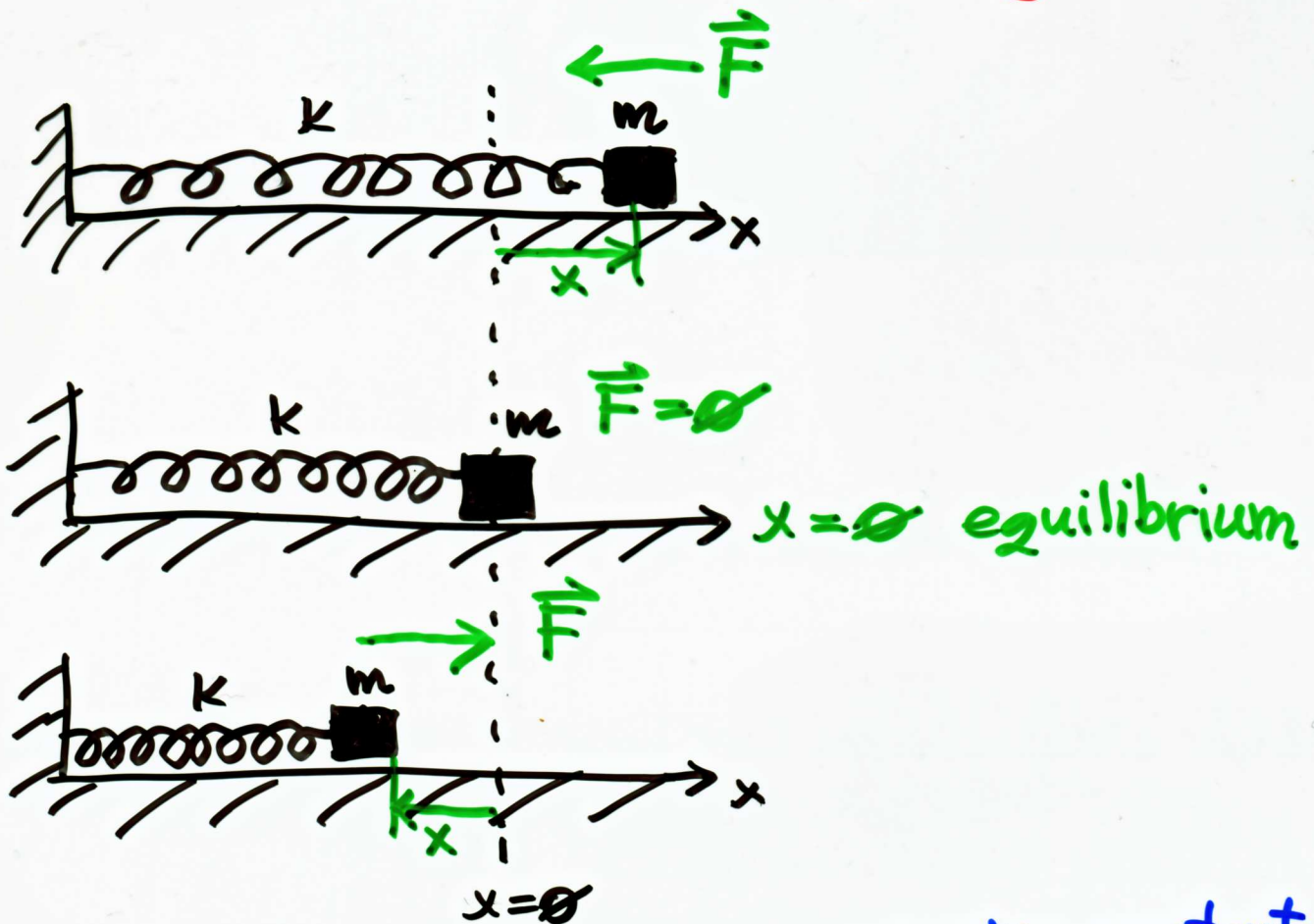
a physical pendulum

at small  $\theta$  ⇒  
simple harmonic motion

(4) damped oscillations

(5) forced oscillations & resonance

# (1) A Particle Attached to a Spring



Hook's Law  $F = -kx$  ← spring constant  
Linear restoring force

How do we mathematically describe the motion?  
Apply the 2<sup>nd</sup> law of Newton

$$F = m \cdot a$$
$$-kx = m \cdot a \quad \& \quad a = d^2x/dt^2$$

12-2

$$-kx = m \frac{d^2x}{dt^2}$$

$$\frac{d^2x}{dt^2} = -\frac{k}{m}x$$

$$\frac{d^2x}{dt^2} + \omega^2 x = 0$$

linear, 2<sup>nd</sup> order  
differential eq.,  
homogeneous

Solution:

$$x(t) = A \cos(\omega t + \phi)$$

A... amplitude of oscillation

$\phi$ ... oscillation phase

$x(t)$  describes position of a particle  
in simple harmonic motion

Check:

$$\frac{dx(t)}{dt} = -A \cdot \omega \sin(\omega t + \phi)$$

$$\left. \begin{aligned} \frac{d^2x(t)}{dt^2} &= -A\omega^2 \cos(\omega t + \phi) \\ \omega^2 x(t) &= +A\omega^2 \cos(\omega t + \phi) \end{aligned} \right\} \oplus \Rightarrow 0$$

Angular frequency  $\omega = \sqrt{\frac{k}{m}}$  [rad/s]

Frequency  $f = \frac{\omega}{2\pi}$  [s<sup>-1</sup>]

Period  $T = \frac{2\pi}{\omega}$

(time needed for one full path of the oscillator)



2π angle (path)  
ω... angular velocity

Phase constant  $\phi$  & amplitude  $A$  in

$$x(t) = A \cos(\omega t + \phi)$$

determined by INITIAL CONDITIONS

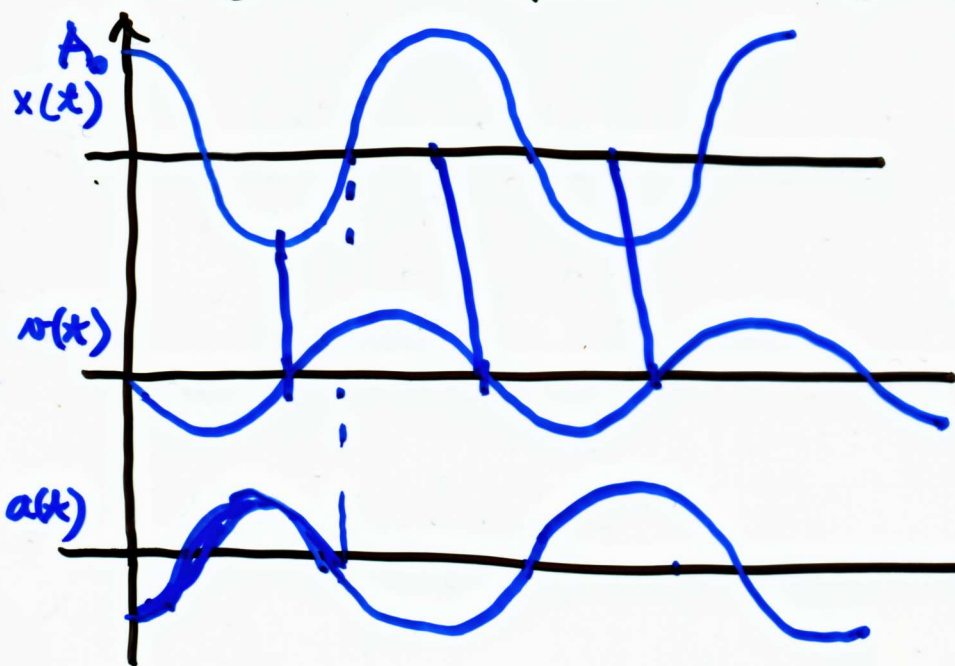
$$x(t=0) = A_0 \Rightarrow A_0 = A \cdot \cos \phi \Rightarrow \phi = 0$$

$A = A_0$

$$x(t) = A_0 \cos(\omega t)$$

$$v(t) = dx(t)/dt = -A_0 \sin(\omega t) \cdot \omega$$

$$a(t) = d^2x(t)/dt^2 = -A_0 \cos(\omega t) \cdot \omega^2$$



$$t=0, \quad t = \frac{1}{4}T$$
$$x_{\max} = A_0, \quad x = 0$$

$$v = 0, \quad v = -A_0 \omega$$

$$a = -\omega^2 A_0, \quad a = 0$$

Quiz: A block on the end of the spring is pulled to position  $x=A$  and released. When is the speed of the block the largest?

~~I (a) at  $x=0$  (equilibrium)~~

I (b) at  $x=-A$  (the left end)

~~8 (c) at  $x=A$  (the right end)~~

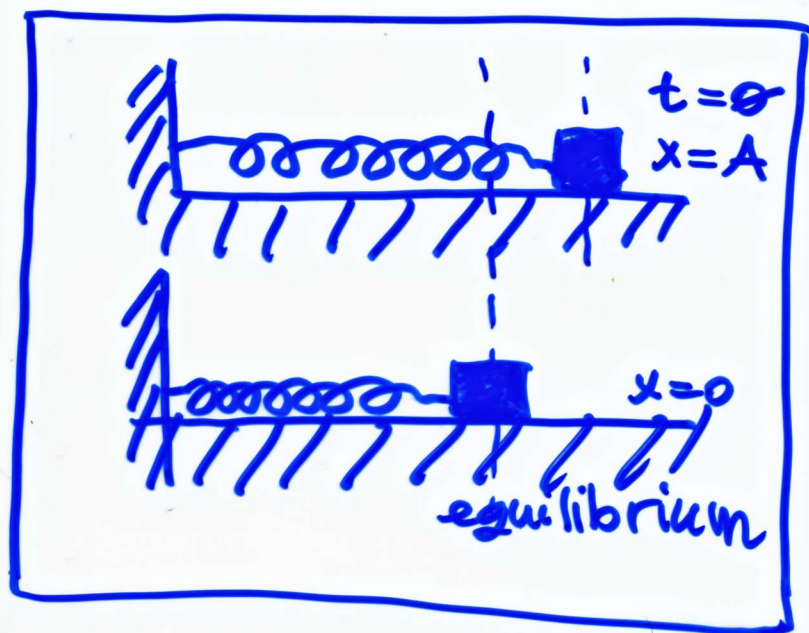
Quiz: What is the total distance the block travels during one full cycle of motion?

~~0 (a)  $A$~~

I (b)  $2A$

2 (c)  $3A$

~~15 (d)  $4A$~~



Example: A particle oscillates as a simple harmonic oscillator along x-axis according to the equation

$$x(t) = (0.05\text{m}) \cdot \cos\left(\pi t + \frac{\pi}{2}\right)$$

(a)  $A = ?$  &  $f = ?$  &  $T = ?$

$$A = 5\text{cm}$$

$$f = \frac{1}{2\pi} \omega = \frac{1}{2\pi} \pi \text{ s}^{-1} = \frac{1}{2} \text{ s}^{-1}$$

$$T = f^{-1} = 2\text{s}$$

(b)  $v = ?$      $a = ?$

$$v = \frac{dx}{dt} = -\pi (0.05\text{m}) \cdot \sin\left(\pi t + \frac{\pi}{2}\right)$$

$$a = \frac{dv}{dt} = -\pi^2 (0.05\text{m}) \cos\left(\pi t + \frac{\pi}{2}\right)$$

(c) initial conditions  $x(t=0) = ?$  &  $v(t=0) = ?$

$$x(t=0) = (5\text{cm}) \cdot \cos\left(\frac{\pi}{2}\right) = 0$$

$$v(t=0) = -\pi \cdot (5\text{cm}) \cdot \sin\left(\frac{\pi}{2}\right) = 1$$

$$= -3.14 \times 5 \frac{\text{cm}}{\text{s}} = \underline{15.7 \text{ cm s}^{-1}}$$



## Energy in Simple Harmonic Oscillator [SHO]

total energy = kinetic e. + potential e.

→ kinetic energy  $K$

$$K = \frac{1}{2} m v^2 = \frac{1}{2} m \omega^2 A^2 \sin^2(\omega t + \phi)$$

→ elastic potential energy  $U$

$$U = \frac{1}{2} k x^2 = \frac{1}{2} k A^2 \cos^2(\omega t + \phi)$$

→ total energy  $E = K + U$

$$E = \frac{1}{2} m \omega^2 A^2 \sin^2(\omega t + \phi) + \frac{1}{2} k A^2 \cos^2(\omega t + \phi)$$

consider  $\omega^2 = k/m \Rightarrow m\omega^2 = k$

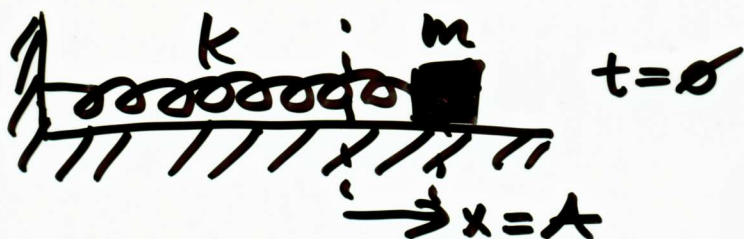
$$E = \frac{1}{2} k A^2 \left[ \sin^2(\omega t + \phi) + \cos^2(\omega t + \phi) \right]$$

||  
1

# Total energy of a SHO

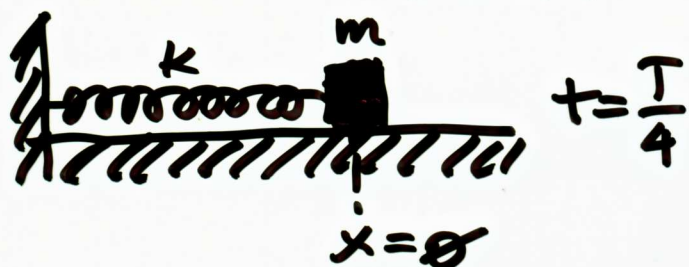
$$E = \frac{1}{2}kA^2 = \text{const.}$$

SHO is conserving the total energy!



$$K = 0$$

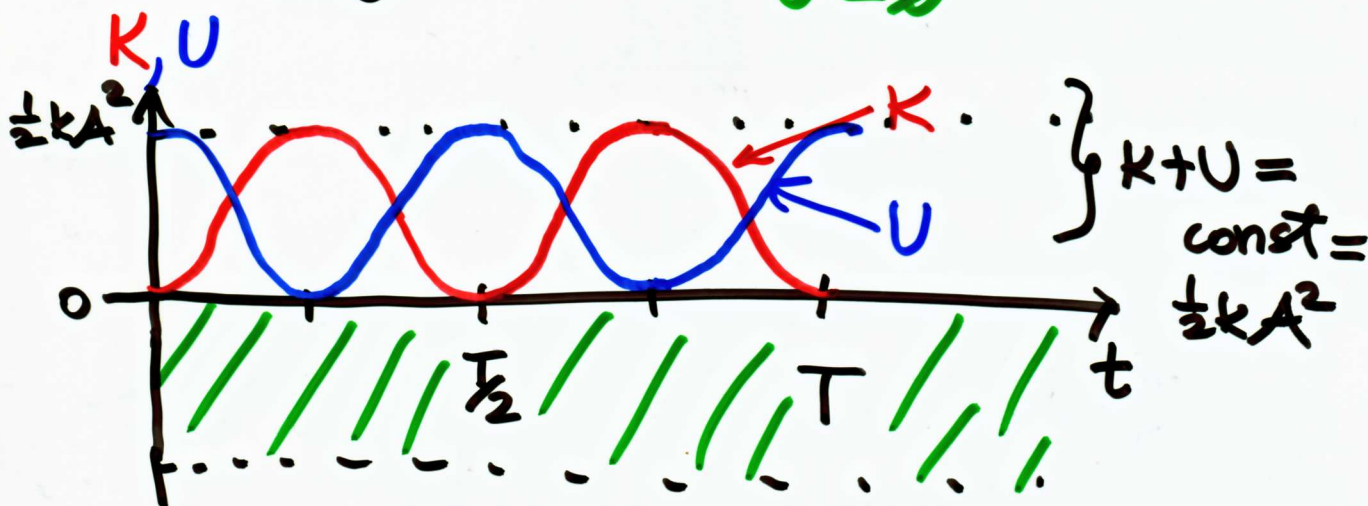
$$U = \frac{1}{2}kA^2 = \text{max}$$



$$K = \frac{1}{2}m\omega^2A^2 =$$

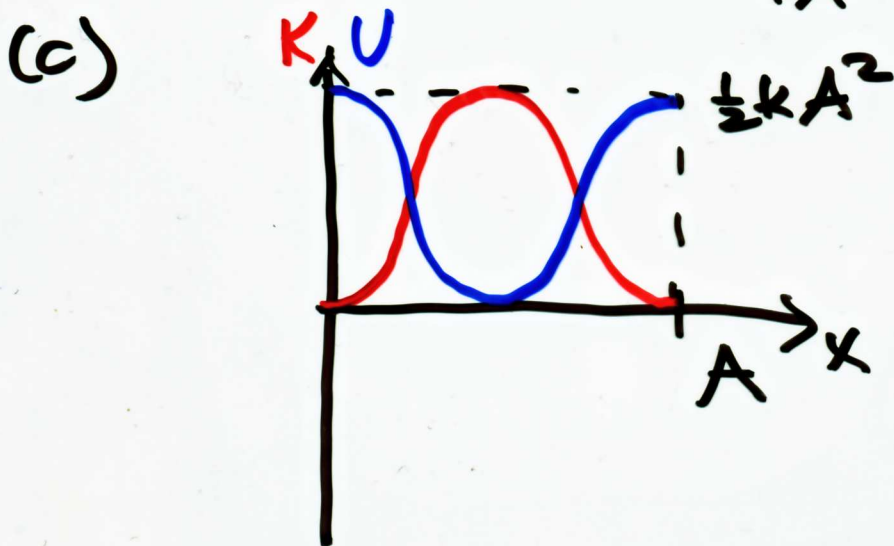
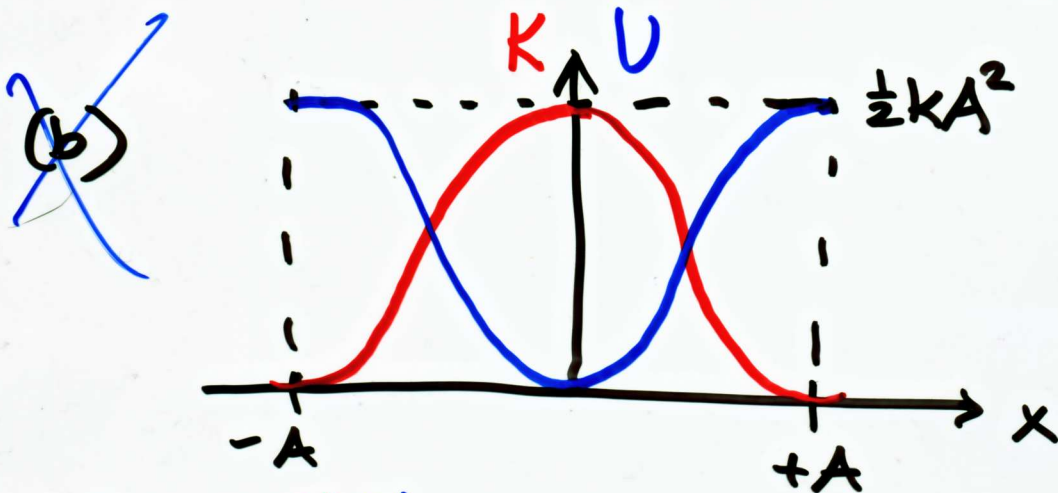
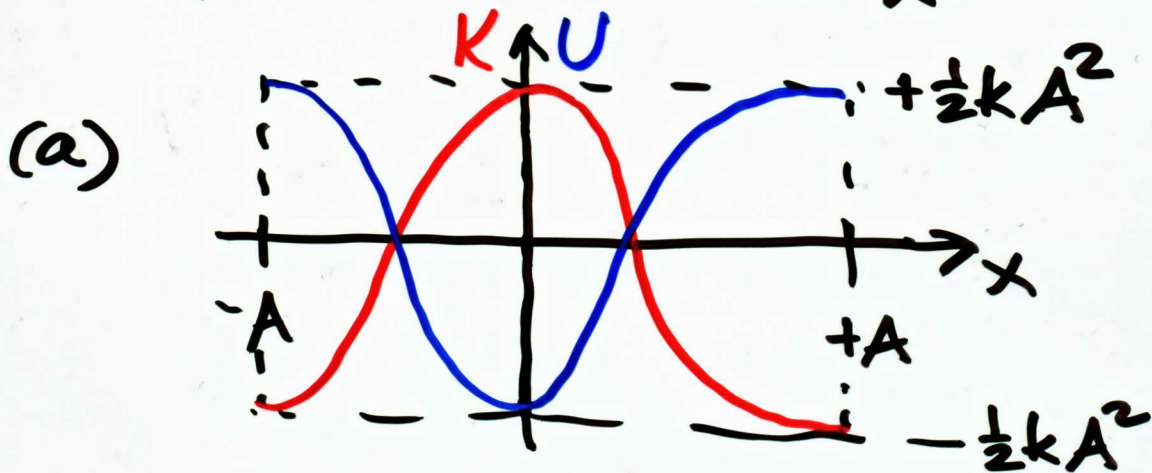
$$= \frac{1}{2}kA^2 = \text{max}$$

$$U = 0$$

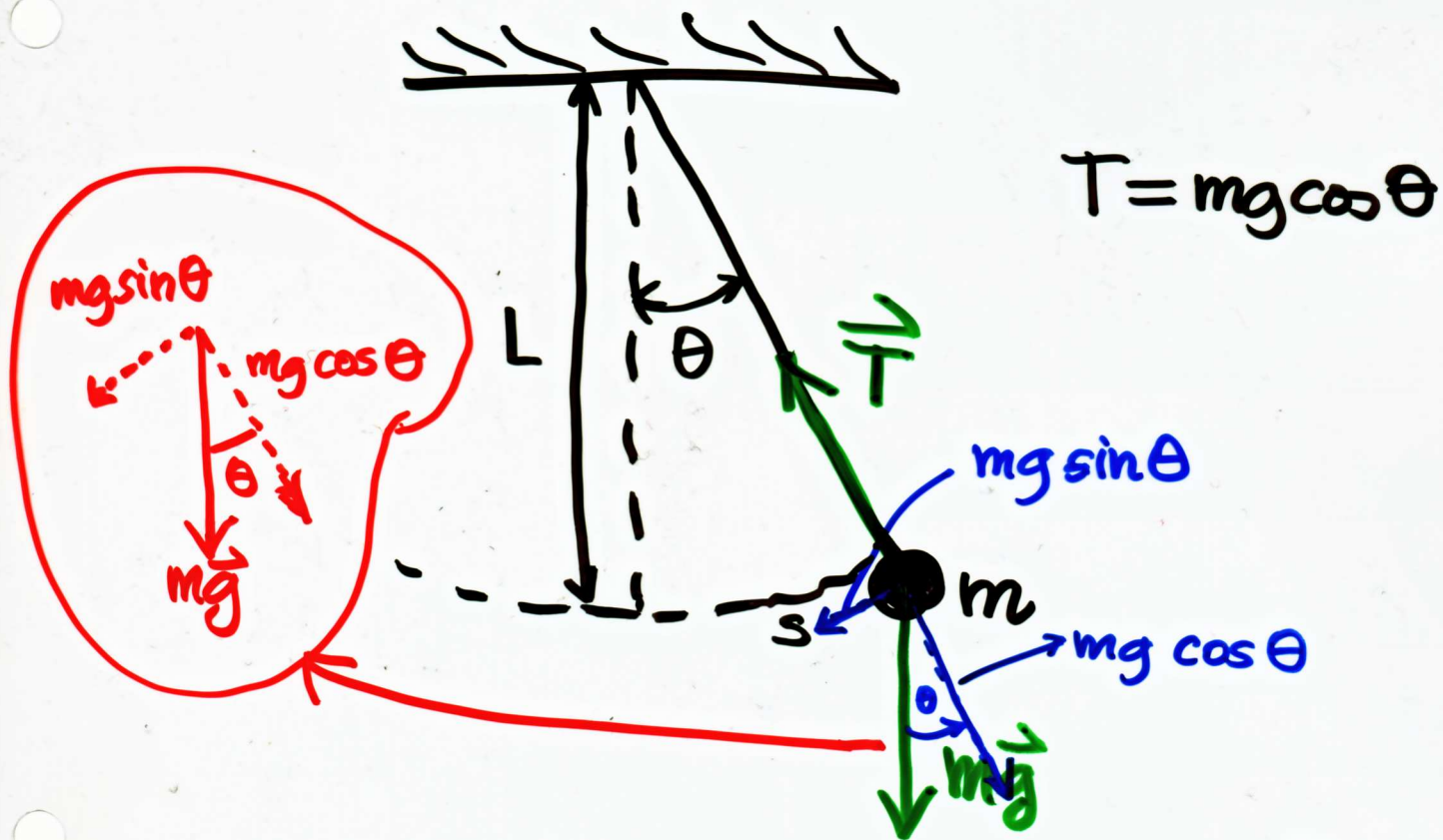


$K$  &  $U$  are oscillating with a period  $T/2$ !  
 $K$  &  $U$  are positive at all times!

Quiz: Consider a block on a "massless" spring. At  $t=0$ ,  $x=A$ , then we let it go. What is the correct dependence of kinetic e.  $K$  & potential e.  $U$  on  $x$ :



## (2) The Simple Pendulum



→ restoring force: along the path  $s$

$$-mg \sin \theta$$

→ Newton 2<sup>nd</sup> Law:  $m \frac{d^2 s}{dt^2} = -mg \sin \theta$

→  $s$  ... circular arc:  $s = L \theta$

$$\frac{d^2 s}{dt^2} = L \cdot \frac{d^2 \theta}{dt^2}$$

$$\Downarrow \quad \cancel{m} L \frac{d^2 \theta}{dt^2} = -\cancel{m} g \sin \theta \quad /: \cancel{m}$$

$$\frac{d^2\theta}{dt^2} = -\frac{g}{L} \cdot \sin\theta$$

Because of  $\sin\theta$  (NOT  $\theta$ ), the simple pendulum does not experience simple harmonic motion in general.

BUT, for small angles  $\theta$ , we can approximate

$\sin\theta \sim \theta$  (where  $\theta$  is measured in radians!) Taylor series

Thus, for  $\theta \ll 1$  : 
$$\frac{d^2\theta}{dt^2} + \frac{g}{L} \theta = 0$$

$$\Rightarrow \theta = \theta_0 \cdot \cos(\omega t + \phi)$$

$$\omega = \sqrt{\frac{g}{L}}$$

$\omega$  does not depend on the mass of the pendulum

$$T = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{L}{g}}$$

All simple pendulums of equal length oscillate with equal periods.

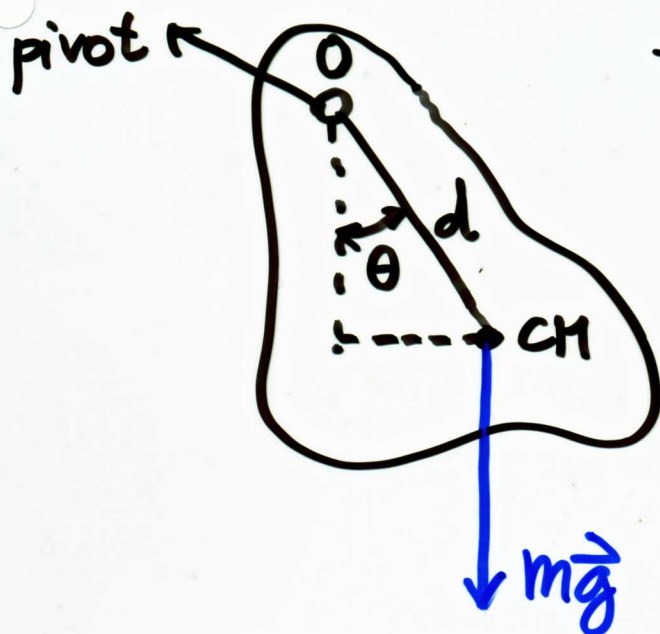
Quiz: Both the block-spring system and the simple pendulum are transported to the Moon. Which of the two will oscillate with a different  $\omega$  on the Moon as compared to the Earth?

(a) <sup>the</sup> block-spring system

(b) ~~the simple pendulum~~

(c) both

### (3) The Physical Pendulum



→ hanging object that oscillates about a fixed axis that does not pass through its CM

→ torque due to gravitational force about the axis O:

$$\vec{\tau} = \vec{r} \times \vec{F}_g \mapsto -mgd \sin \theta$$

→ Newton's 2<sup>nd</sup> Law:

$$\sum \vec{\tau} = I \vec{\alpha}, \quad I \text{ moment of inertia}$$

$$\vec{\alpha} = \frac{d^2\theta}{dt^2} \text{ angular acceleration}$$



$$I \frac{d^2\theta}{dt^2} = -mgd \sin \theta$$

→ small angle approximation

$$\frac{d^2\theta}{dt^2} = - \frac{mgd}{I} \theta$$

Angular frequency  $\omega$  & period  $T$

$$\left. \begin{aligned} \omega &= \sqrt{\frac{mgd}{I}} \\ T &= 2\pi \sqrt{\frac{I}{mgd}} \end{aligned} \right\} *$$

The physical pendulum is a generalized pendulum. You can use the above Eqs. \* to find the simple pendulum  $\omega$  &  $T$ .

For simple pendulum:

$$I = md^2 \Rightarrow \omega = \sqrt{\frac{g}{d}}$$

$$T = 2\pi \sqrt{\frac{d}{g}}$$



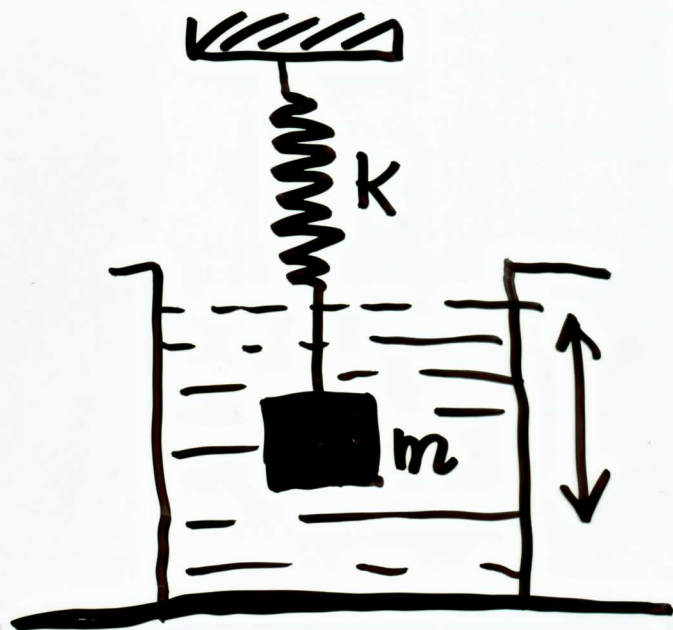
## (4) Damped Oscillations

SHO was considered to occur in ideal frictionless conditions.

In a more realistic description, we need to take into account FRICTION, a resistive force which opposes motion



Total energy of a damped oscillator decreases with time.



Example of a damped oscillator

# Mathematical description

→ linear resistive force:  $-b\vec{v} = -b \frac{d\vec{x}}{dt}$

→ Hooke's force:  $-kx$

→ Newton's 2<sup>nd</sup> law:  $\sum \vec{F} = m\vec{a}$

$$m \frac{d^2x}{dt^2} = -kx - b \frac{dx}{dt} \quad (1D)$$

$$\frac{d^2x}{dt^2} + \frac{b}{m} \frac{dx}{dt} + \omega_0^2 x = 0$$

$$\omega_0 = \sqrt{\frac{k}{m}}$$

underdamped  
osc.

$$b < \sqrt{4mk}$$

critically  
damped  
osc.

$$b = \sqrt{4mk}$$

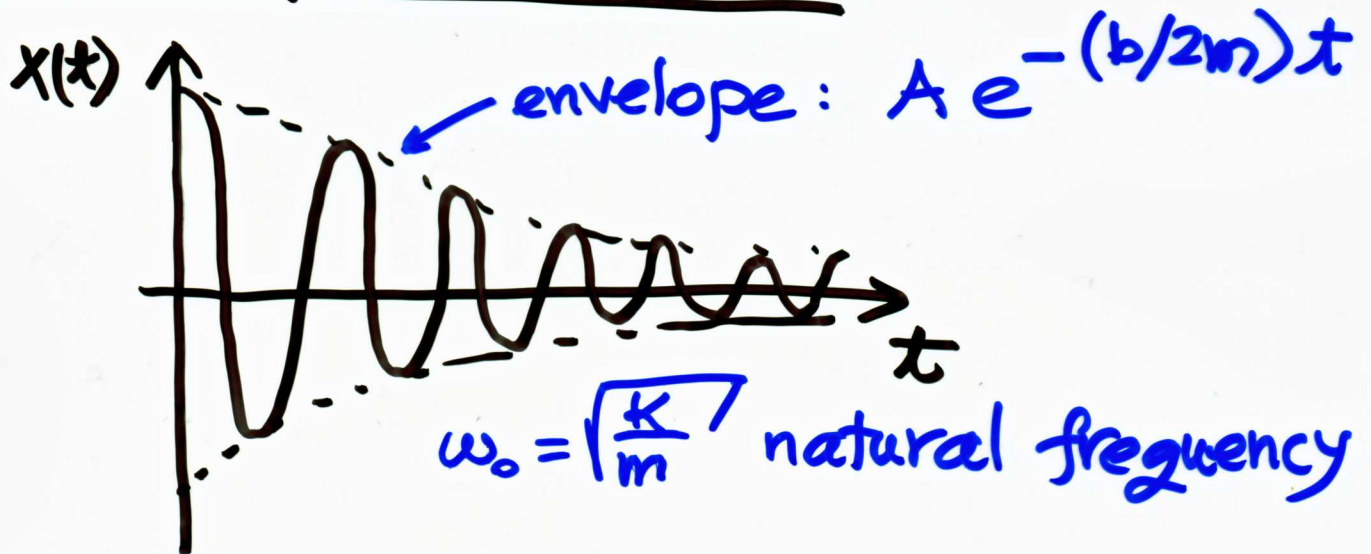
overdamped  
osc.

$$b > \sqrt{4mk}$$

$$x(t) = A_0 e^{-(b/2m)t} \cos(\omega t + \phi)$$

$$\omega = \sqrt{\omega_0^2 - \left(\frac{b}{2m}\right)^2}$$

## Underdamped oscillator:



→ envelope exponential decay  
(the larger the  $b$ , the faster the decay)

→ frequency  $\omega$  DEPENDS on  $b$ :

$$\omega = \sqrt{\omega_0^2 - \left(\frac{b}{2m}\right)^2}$$

⇓

Oscillations slower is a more  
viscous liquid (larger  $b$ )

critically damped & overdamped  
⇒ no oscillations,  
only exponential decay  
of  $x(t)$

Example: Calculate how the total  $E$  of a damped oscillator decreases with time.  
Assume that at  $t=0$ ,  $x=A$ .

## (5) Forced Oscillations & Resonance

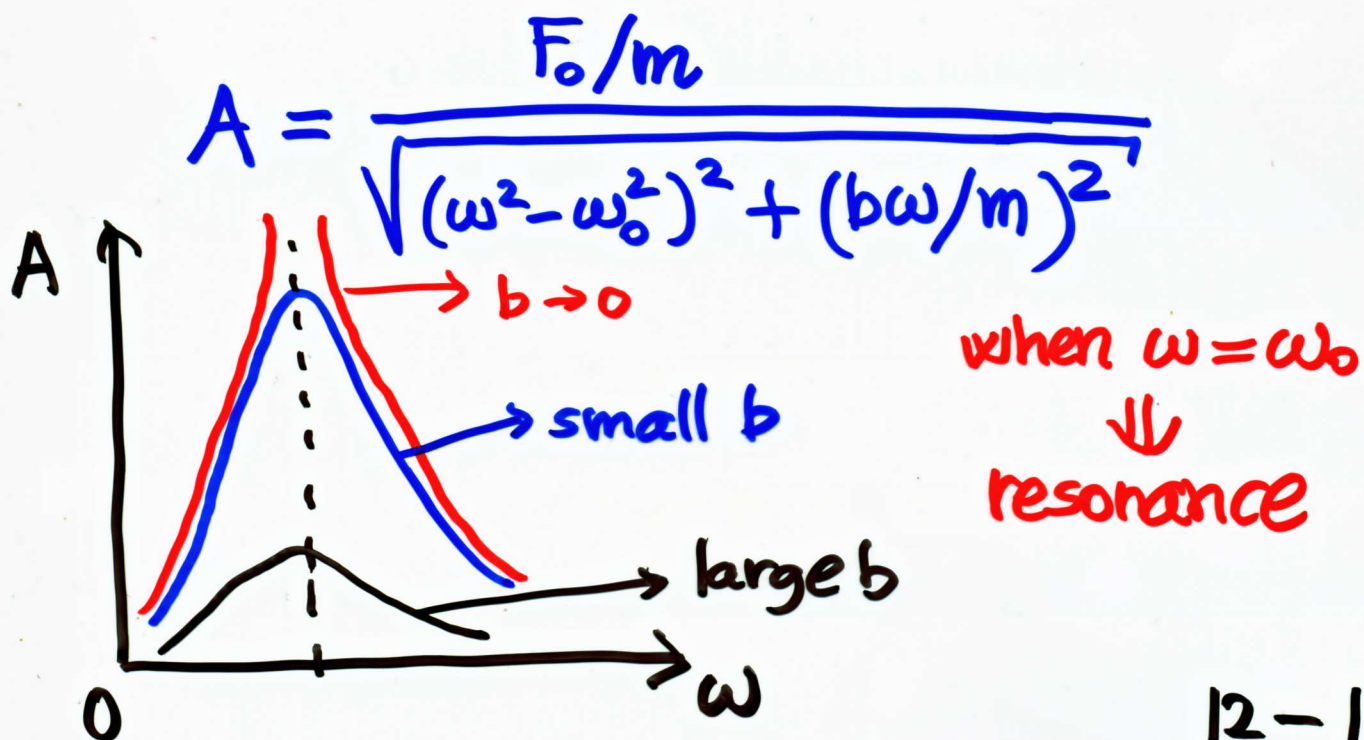
→ to avoid loss of mechanical E with time in a damped oscillator, apply an external force to provide positive work on the system:

$$F = F_0 \sin \omega t$$

→ after some initial time, the system driven by the external force gets into a steady state:

$$x(t) = A \cos(\omega t + \phi)$$

where:



# Ch. 13: Mechanical Waves

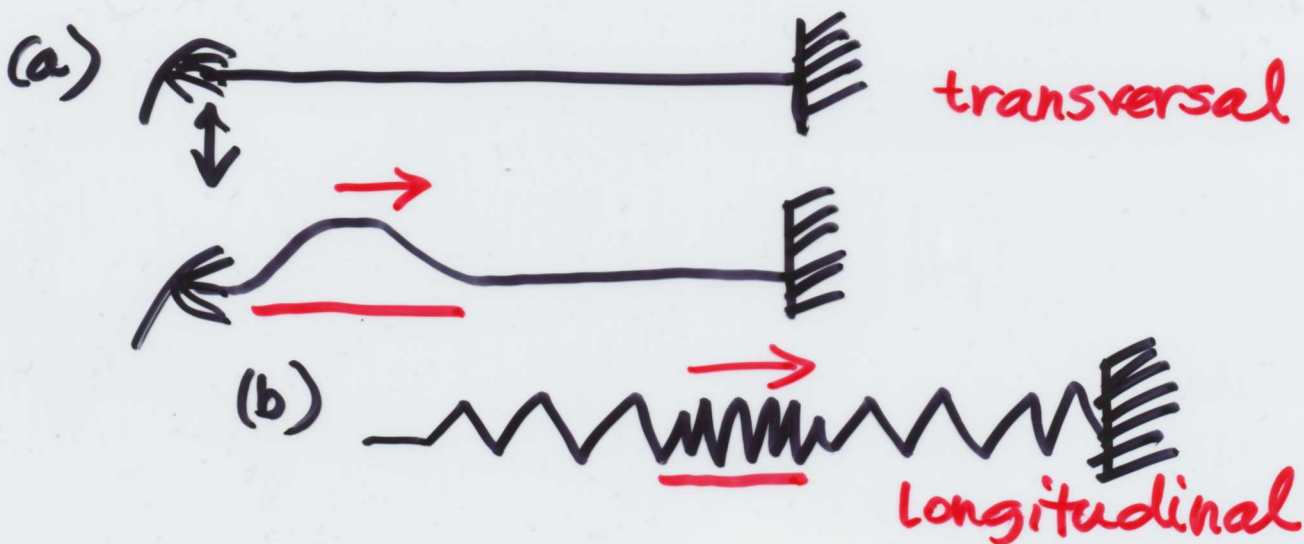
a pebble dropped into a pond  
⇒ disturbance ⇒ water waves

Mechanical waves:

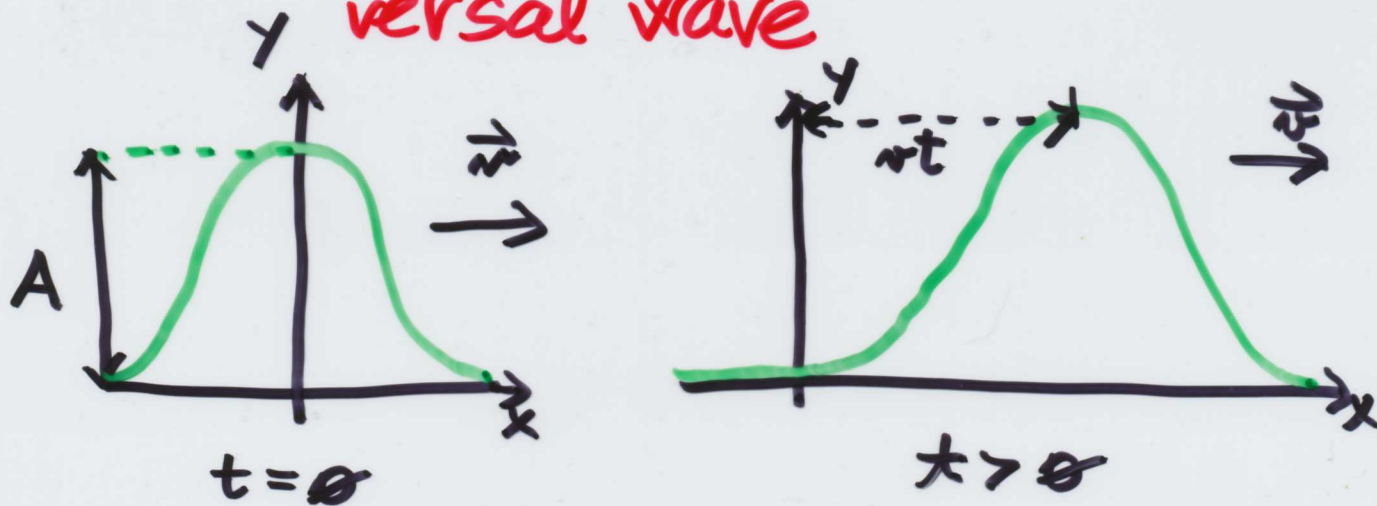
- (a) → transversal (metal string)
- (b) → longitudinal (sound)

Requirements for creating the mechanical waves

- (1) source of disturbance
- (2) medium that can be disturbed
- (3) physical mechanism through which elements of media influence each other



# Mathematical description of a trans- versal wave



→ consider a wave that does not change the shape with time  
⇒ no dispersion

→ mathematical formulation:

$$y(x, t) = y(x - vt, 0)$$

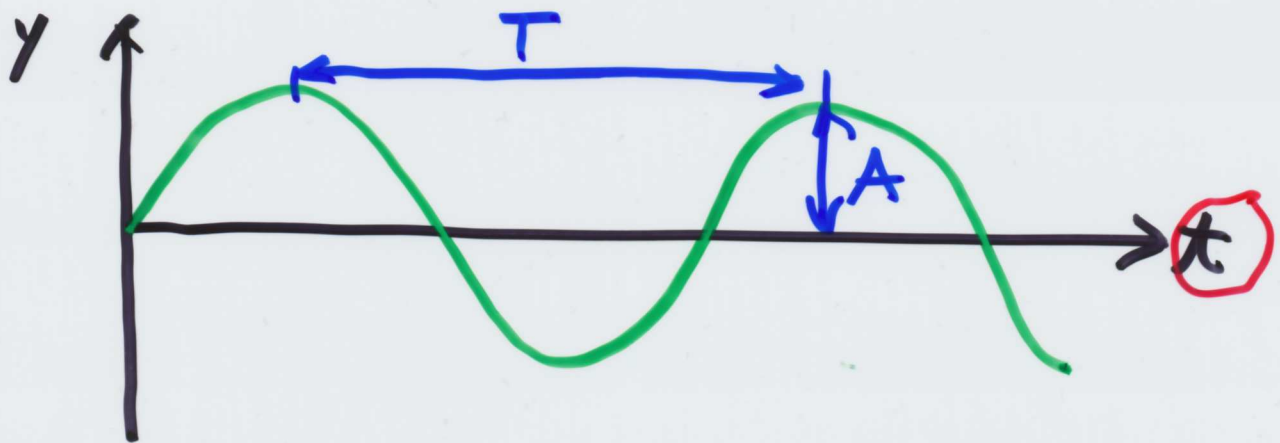
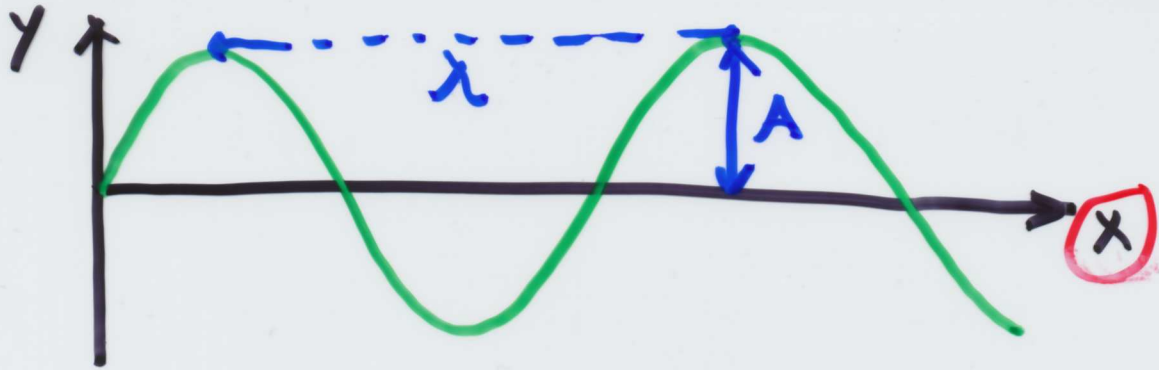
for a wave moving  $L \rightarrow R$

$$y(x, t) = y(x + vt, 0)$$

for a wave moving  $R \rightarrow L$

→  $y(x, t)$  represents the y-position of any element of the string located at  $x$  at time  $t$

# Wave Model $y = y(x, t)$



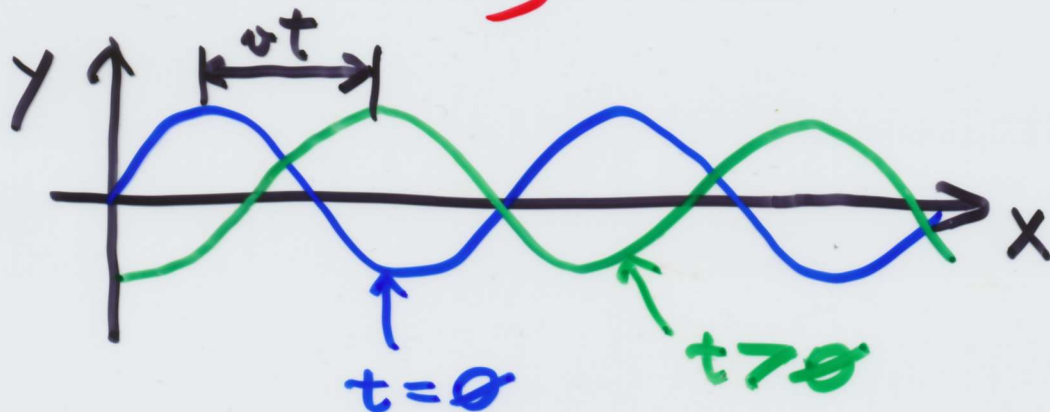
$\lambda$ ... wavelength  
 $A$ ... amplitude  
 $T$ ... period

$$T = \frac{1}{f}$$
$$\omega = 2\pi f$$

Each element (of a string/water) oscillates vertically in the  $y$  direction with simple harmonic motion.



# The Travelling Wave



$$\rightarrow t=0 : y = A \sin\left(\frac{2\pi}{\lambda}x\right)$$

$$\rightarrow t>0 : y = A \sin\left[\frac{2\pi}{\lambda}(x - vt)\right]$$

$$v = \frac{\lambda}{T} \quad (T \dots \text{period})$$

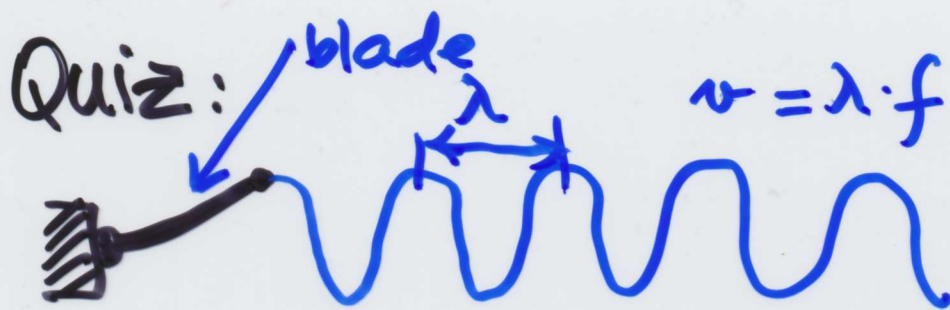
Most common form of a travelling wave :

$$y(x,t) = A \cdot \sin(kx - \omega t)$$

$$k = \frac{2\pi}{\lambda} \quad \text{wave number}$$

$$\omega = \frac{2\pi}{T} = 2\pi f \quad \text{angular frequency}$$

$$v = \frac{\omega}{k} \quad \text{OR} \quad v = \lambda f = \frac{\lambda}{T}$$



A sinusoidal wave with frequency  $f$  is traveling along a stretched string. Then, we replace the blade such that the new one is oscillating with  $2f$ , producing a wave with frequency  $2f$ , traveling on the same string.

What is the wavelength of the new wave?

(a)  $\lambda_N = 2\lambda$

(b)  $\lambda_N = \lambda$

(c)  $\lambda_N = \frac{1}{2}\lambda$

(d) impossible to determine

# The Linear Wave Equation

Consider transverse traveling wave

$$y(x, t) = A \cdot \sin(kx - \omega t)$$

Transverse velocity & acceleration

$$\rightarrow v_y = \frac{\partial y}{\partial t} = \left. \frac{dy}{dt} \right|_{x=\text{const}} = -\omega A \cos(kx - \omega t)$$

$$\rightarrow a_y = \frac{\partial^2 y}{\partial t^2} = \left. \frac{d^2 y}{dt^2} \right|_{x=\text{const}} = -\omega^2 A \cdot \sin(kx - \omega t)$$

are periodic functions, out of phase

We obtained:  $\frac{\partial^2 y}{\partial t^2} = -\omega^2 y(x, t)$

Second derivative with respect to  $x$ ,

$$\frac{\partial^2 y}{\partial x^2}: \quad \frac{\partial^2 y}{\partial x^2} = -k^2 y(x, t)$$

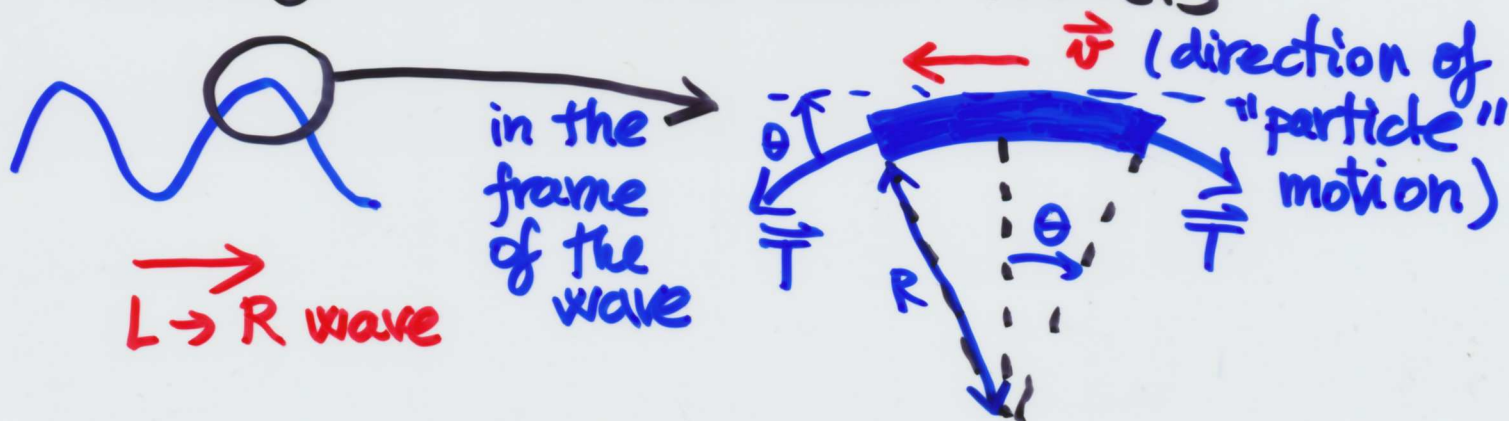
Take into account  $v = \omega/k$



$$\frac{\partial^2 y}{\partial x^2} = -\frac{k^2}{\omega^2} \cdot \omega^2 y(x, t) = +\frac{1}{v^2} \frac{\partial^2 y}{\partial t^2}$$

# The Speed of Transverse Waves on Strings

Important: the wave speed depends only on the properties of the medium through which the wave travels



radial force

$$F_T = 2T \sin \theta \approx 2T \theta \quad (\text{SAA})$$

$$\left. \begin{aligned} m &= (\mu \cdot \Delta s) \\ \Delta s &= R(2\theta) \end{aligned} \right\} \Rightarrow$$

$\mu$  [ $\text{g m}^{-1}$ ]  
mass/unit length

$$\Rightarrow m = 2\mu R \theta$$

$$F_T = \frac{mv^2}{R}$$

centripetal force  
(Newton's 2<sup>nd</sup> law)

$$2T \theta = \frac{2\mu R \theta v^2}{R} \Rightarrow$$

$$v = \sqrt{\frac{T}{\mu}}$$

# The Rate of Energy Transfer



$$\left. \begin{aligned} \Delta K &= \frac{1}{2} (\Delta m) v_y^2 = \\ &= \frac{1}{2} (\mu \Delta x) v_y^2 \end{aligned} \right\} dK = \frac{1}{2} \mu \omega^2 A^2 \cos^2(kx - \omega t)$$

$$\left. \begin{aligned} K_\lambda &= \int_0^\lambda dK = \frac{1}{4} (\mu \omega^2 A^2 \lambda) \\ U_\lambda &= K_\lambda \end{aligned} \right\} E_\lambda = \frac{1}{2} \mu \omega^2 A^2 \lambda$$

$$P = \frac{E_\lambda}{T} = \frac{1}{2} \mu \omega^2 A^2 \left( \frac{\lambda}{T} \right) = v$$

$$P = \frac{1}{2} \mu \omega^2 A^2 v$$

The rate of energy transfer  $\propto$   
square of the angular frequency  $\omega$   
and the square of amplitude  $A$ !

Quiz: How would you increase the rate of energy transfer THE MOST?

(a) reduce the linear mass density  $\mu$

(b) doubling  $\lambda$

(c) doubling  $T$  (tension in the string)

~~(d) doubling the amplitude~~