## INTRODUCTION TO RELATIVITY Winter 2018-2019 Prof. Michael S. Vogeley Homework Assignment 4 Solutions

Spacetime ch.L, problems 6 (a only), 7, 8

## 6. Transformation of Angles

(a) In the rocket frame, the meter stick lies at angle  $\phi'$  with respect to the x' axis. In the rocket frame  $\Delta x' = (1 \text{ m}) \cos \phi'$  and  $\Delta y' = (1 \text{ m}) \sin \phi'$ . In the laboratory frame, the stick appears contracted along the direction of travel, thus  $\Delta x = \Delta x'/\gamma = (1 \text{ m}) \cos \phi'/\gamma$ . There is no change tranverse to the direction of travel,  $\Delta y = \Delta y' = (1 \text{ m}) \sin \phi'$ . The angle with the lab frame x axis is larger than in the rocket frame,  $\tan \phi = \Delta y/\Delta x = \gamma \Delta y'/\Delta x'$ , so  $\phi = \tan^{-1}(\gamma \tan \phi')$ . The length of the meter stick in the lab frame d is given by  $d^2 = (\Delta x)^2 + (\Delta y)^2 = [(1 \text{ m}) \cos \phi'/\gamma]^2 + [(1 \text{ m}) \sin \phi']^2$ . A little algebra yields  $d = (1 \text{ m})\sqrt{1 - v^2 \cos^2 \phi'}$ .

## 7. Transformation of y-velocity

In the rocket frame, the particle moves at  $v'_y = \Delta y'/\Delta t'$ . Assume that the rocket moves at  $v_{rel}$  in the x direction relative to the lab frame. Apply the Lorentz transformation: No effect on distances in the transverse dimension, so in the lab frame  $\Delta y = \Delta y'$ . But we do have to transform the time,  $\Delta t = \gamma \Delta t'$ . Thus, the y-velocity in the lab frame is  $v_y = \Delta y/\Delta t = \Delta y'/(\gamma \Delta t') = \sqrt{1 - v_{rel}^2}v'_y$ . In the rocket frame, the particle has  $v'_x = 0$ . But in the lab frame, it moves  $\Delta x = v_{rel}\Delta t$ , thus  $v_x = v_{rel}$ .

## 8. Transformation of velocity direction

As suggested in the problem, solve this by transforming the space and time intervals. In the rocket frame, in time  $\Delta t'$  the particle moves in the y' direction an amount

$$\Delta y' = v' \sin \phi' \Delta t'$$

and moves in the x' direction an amount

$$\Delta x' = v' \cos \phi' \Delta t'$$

Now transform  $\Delta y'$  and  $\Delta x'$  into the lab frame: There is no change in the transverse displacement,

$$\Delta y = \Delta y'$$

but the displacement along the direction of travel and the time does change:

$$\Delta x = \gamma \Delta x' + \gamma v_{rel} \Delta t'$$

The tangent of the angle in the lab frame is then

$$\tan \phi = \frac{\Delta y}{\Delta x} = \frac{\Delta y'}{\gamma \Delta x' + \gamma v_{rel} \Delta t'}$$

Plug in the expressions for  $\Delta y'$  and  $\Delta x'$  in the rocket frame from above,

$$\tan \phi = \frac{v' \sin \phi' \Delta t'}{\gamma(v' \cos \phi' \Delta t') + \gamma v_{rel} \Delta t'}$$

The  $\Delta t'$  cancels in numerator and denominator and you can also factor out v' to yield

$$\phi = \tan^{-1} \left( \frac{\sin \phi'}{\gamma(\cos \phi' + v_{rel}/v')} \right)$$

This problem differs from problem 6 because this involves the motion of a particle, rather than the appearance of an object. In the limit of very large  $v_{rel}$ , the argument of  $\tan^{-1}$  goes to zero, thus  $\phi \to 0$ . This makes sense because the x displacement becomes huge compared to the y displacement in fixed time. In problem 6, as  $v_{rel} \to \infty$ , the argument of  $\tan^{-1}$  becomes large, thus  $\phi \to 90^{\circ}$