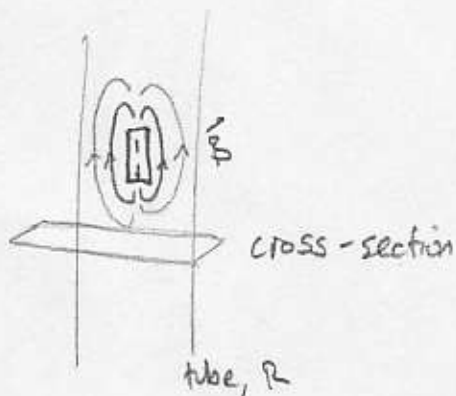


PHYS 115] HW #6 solutions

#1]

Consider a cross-section of the tube



As the magnet falls - the flux $\Phi = \int \vec{B} \cdot d\vec{A}$ increases,

this creates an emf to oppose the change in flux, This emf creates its own magnetic field, since

in our cross-section we now have a current $I = \epsilon_{ind}/R$. The \vec{B} field here will slow down the magnet.

#2]

$$|\vec{B}(t)| = B_0 + bt^3 \quad \frac{d}{dt} |\vec{B}| = 3bt^2$$

Inside $r < R$



$$|\mathcal{E}| = \frac{d\Phi}{dt} = -\frac{d}{dt} \int \vec{B} \cdot \hat{n} dA = -\frac{d}{dt} \pi r^2 |\vec{B}| = -3bt\pi r^2 t^2$$

$$E = 3/2 b r t^2$$

At a larger radius $r_2 > R$ only the portion from $r \leq R < r_2$ counts,

hence $|\mathcal{E}| = 3b\pi R^2 t^2$

$$\mathcal{E} = \int \vec{E} \cdot d\vec{l} = E 2\pi r_2$$

$$E = 3/2 b R^2 t^2 / r_2$$

Since the flux is growing out of the page, the emf must go clockwise



$t=0$



$0 < t'$

#3]

$$B = .07 + .03t^2$$

$$\frac{dB}{dt} = .06t$$



$$A_{sol} = \pi r^2 ; r = .03m$$

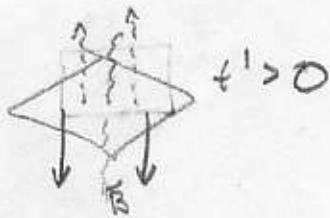
$$|\mathcal{E}_{ind}| = \frac{d\Phi}{dt} = \oint \vec{B} \cdot \hat{n} dA = \frac{d}{dt} N \left[\int_{\text{inside solenoid}} \vec{B} \cdot \hat{n} dA + \int_{\text{outside solenoid}} \vec{B} \cdot \hat{n} dA \right]$$

$$= N \cdot \frac{d}{dt} [B\pi r^2] = N\pi r^2 (.06)t$$

$$I(t=2) = \frac{\mathcal{E}(t=2)}{R} = .0136A$$



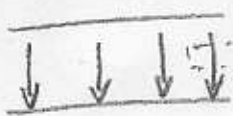
t=0



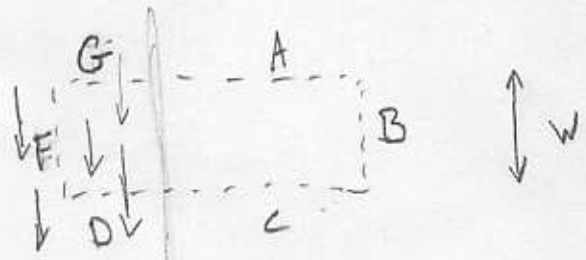
t' > 0

Current must run clockwise in the picture shown

#4]



take path as shown

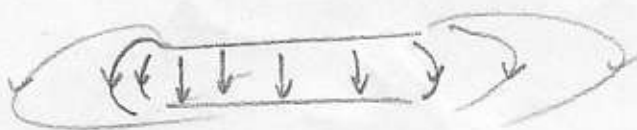


$$\oint \vec{E} \cdot d\vec{l} = 0 \quad \text{for } A, B, C \text{ since } \vec{E} = 0$$

$$\oint \vec{E} \cdot d\vec{l} = 0 \quad \text{for } G, D \text{ since } \vec{E} \cdot d\vec{l} = 0$$

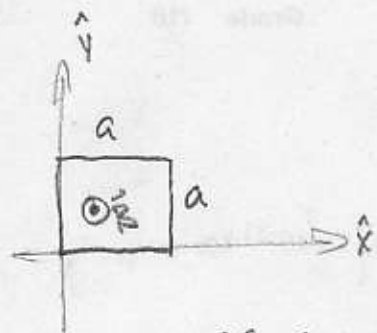
$$\oint \vec{E} \cdot d\vec{l} = Ew \quad \text{for path } F \text{ but } \oint \vec{E} \cdot d\vec{l} = 0 \neq Ew !$$

Hence, this must not be the field.



A more realistic field — fringes

#5]



$$\frac{dB}{dt} = \frac{d}{dt} 4t^2y = 8ty$$

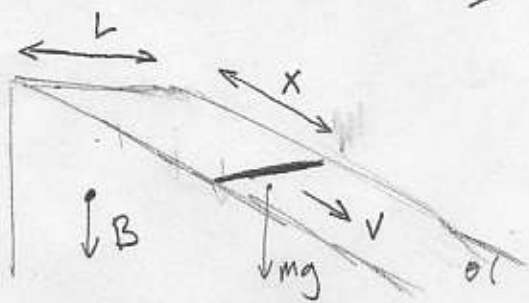
$$\Phi = \int \vec{B} \cdot \hat{n} dA$$

$$= \int_0^a \int_0^a \vec{B}(x,y) \cdot \hat{n} dx dy = \int_0^a \int_0^a 4t^2y dx dy$$

$$= (4t^2a) \left. \frac{y^2}{2} \right|_0^a = 2t^2a^3$$

Hence $|\mathcal{E}| = \frac{d\Phi}{dt} = 4t^2a^3$ $\mathcal{E}(t=2s) = 8 \cdot 10^{-5} \text{ V}$, clockwise

#6]



IF v is in the direction down the plane then

$$F_g = mg \sin \theta$$

$$\Phi = \int \vec{B} \cdot \hat{n} dA = (\vec{B} \cdot \hat{n}) Lx = B \cos \theta Lx$$

$$\frac{d\Phi}{dt} = B \cos \theta L \frac{dx}{dt} = BLv \cos \theta$$

$$I_{\text{ind}} = \frac{d\Phi}{dt} / R = \frac{BLv \cos \theta}{R}$$

The force on a current in a B field is $F_B = I \vec{L} \times \vec{B} = ILB$

* but * this force is not along the direction of v , $F_{B, \text{along } v} = ILB \cos \theta$

Steady state implies $F_g = F_B$, $\frac{BL^2 \cos^2 \theta}{R} v = mg \sin \theta$

$$v = \frac{mg \sin \theta}{L^2 B^2 \cos^2 \theta} R$$

rate of potential energy lost

$$= \frac{d}{dt} U = \frac{d}{dt} mg \sin \theta x = mg \sin \theta v = mg \sin \theta \left(\frac{mg \sin \theta}{L^2 B^2 \cos^2 \theta} R \right)$$

same!

$$P = \mathcal{E}^2 / R = \frac{(d\Phi/dt)^2}{R} = \frac{(BL \cos \theta)^2 v^2}{R} = \frac{(BL \cos \theta)^2}{R} R^2 \frac{(mg \sin \theta)^2}{(L^2 B^2 \cos^2 \theta)^2} = \frac{(mg \sin \theta)^2 R}{L^2 B^2 \cos \theta}$$

effect would still hold if $\vec{B} \rightarrow -\vec{B}$ as the flux would be the same