Lecture 6: Entropy Rules! (contd.)

Lecturer:

Brigita Urbanc

Office: 12-909

(E-mail: brigita@drexel.edu)

Course website:

www.physics.drexel.edu/~brigita/COURSES/BIOPHYS_2011-2012/

Boltzmann distribution (by counting: Textbook, pages 230-231)

Information theory:

based on the constraints (for example the total energy) derive the least biased probability distribution

Shannon entropy:

$$S(p_1, p_2, p_3, ..., p_N) = -\sum_{i=1}^{N} p_i \ln p_i$$

where \mathbf{p}_{i} is the probability of the system to be in the i-th microstate.

For example, if nothing is known about the system except that there is N microstates, then for all $p_i=1/N$ and $S=ln\ N$ (the maximal value).

How do we formally derive this result?

Maximize Shannon entropy S' using Lagrange multiplier method for each constraint:

$$S' = -\sum_{i=1}^{N} p_i \ln p_i - \lambda \left[\sum_{i=1}^{N} p_i - 1 \right]$$

Maximization equations:

$$\frac{\partial S'}{\partial \lambda} = 0 \longrightarrow -\sum_{i=1}^{N} p_i + 1 = 0$$

$$\frac{\partial S'}{\partial p_i} = 0 \longrightarrow -\ln p_i - 1 - \lambda = 0$$

$$p_i = \exp(-1 - \lambda)$$

Note that the probabilities **p**; do not depend on i!

$$\sum_{i=1}^{N} p_i = 1 \rightarrow \sum_{i=1}^{N} \exp(-1-\lambda) = 1 \rightarrow \exp(-1-\lambda) = 1/N$$

$$p_i = 1/N$$

Boltzmann distribution is a maximum entropy distribution with a fixed average energy:

$$S' = -\sum_{i} p_{i} \ln p_{i} - \lambda \left[\sum_{i} p_{i} - 1 \right] - \beta \left[\sum_{i} p_{i} E_{i} - \langle E \rangle \right]$$

$$-\ln p_i - 1 - \lambda - \beta E_i = 0 \rightarrow p_i = \exp(-1 - \lambda) \exp(-\beta E_i)$$
$$\sum_i p_i = 1 \rightarrow \exp(1 + \lambda) = \sum_i \exp(-\beta E_i) = Z$$

$$p_i = \exp(-\beta E_i) / \sum_i \exp(-\beta E_i)$$

Ideal gas approximation: the interaction range is small in comparison to the mean spacing between molecules

For a system with N independent variables $(X_1, X_2, ..., X_N)$ the probability distribution can be factorized:

$$P(x_1, x_2, ... x_N) = P(x_1) P(x_2) ... P(x_N)$$

 uniform spatial distribution of ideal gas molecules.

Microscopic state of the system is described by (x, p_x) so the key information that we need to derive is $P(p_x)$ knowing the kinetic energy of $E = p_x^2/(2m)$.

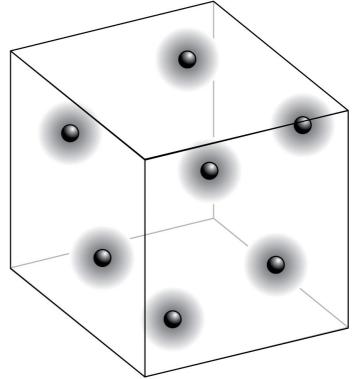


Figure 6.19 Physical Biology of the Cell (© Garland Science 2009)

Probability distribution for a system with average energy is:

$$\exp[-\beta p_x^2/(2m)]$$

$$P(p_y) = \frac{-}{}$$

$$\sum_{\text{states}} \exp[-\beta p_x^2/(2m)]$$

Instead of a sum, we use an integral over a continuous variable $\mathbf{p}_{\mathbf{y}}$:

$$\sum_{\text{states}} \rightarrow \int_{-\infty}^{\infty} dp_{x}$$

where we will use the integral:

$$\int_{-\infty}^{\infty} \exp(-\alpha p_x^2) dp_x = \sqrt{\pi/\alpha}$$

According to the equipartition theorem, each degree of freedom is associated with:

$$\langle \mathbf{E} \rangle = \frac{1}{2} \mathbf{k}_{\mathbf{B}} \mathbf{T}$$

Calculate $\langle E \rangle$ using the Boltzmann distribution:

$$Z = \int_{-\infty}^{\infty} \exp[-\beta p_x^2/(2m)] dp_x = \sqrt{2m\pi/\beta}$$

$$\langle E \rangle = \int_{-\infty}^{\infty} p_x^2/(2m) \exp[-\beta p_x^2/(2m)] dp_x$$

We use the following trick:

$$\langle \mathbf{E} \rangle = - \mathbf{Z}^{-1} \partial \mathbf{Z} / \partial \boldsymbol{\beta} = \frac{1}{2} \boldsymbol{\beta}^{-1}$$

so we showed that the Lagrange multiplier $\beta = (k_B^T)^{-1}$.

Free energy and chemical potential of a *dilute solution*: Application of a lattice model (and ideal gas approx.)

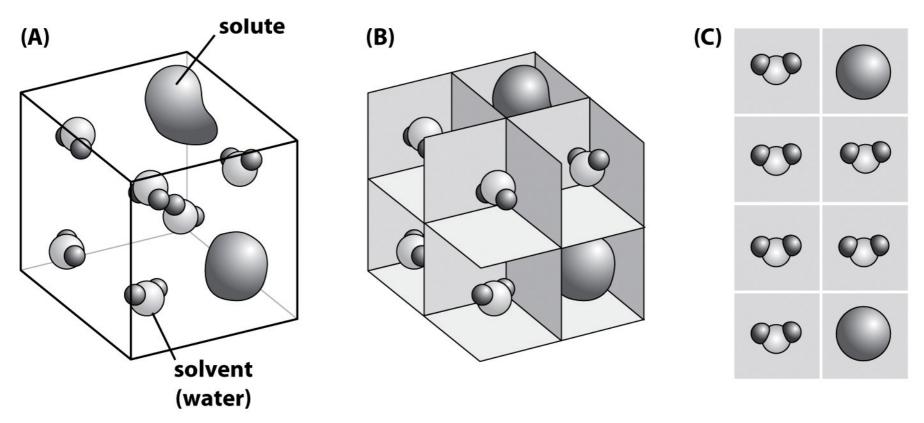


Figure 6.20 Physical Biology of the Cell (© Garland Science 2009)

Configurational entropy of N objects placed into W available spots:

$$\Omega$$
 !

$$W(N,\Omega) = \frac{1}{S} = k_B \ln W$$

 $N! (\Omega - N)!$

How to calculate the chemical potential of a dilute solutions?

$$\begin{split} & \mu_{\text{SOLUE}} &= \left(\partial G_{\text{TOT}} / \partial N_{\text{S}}\right)_{\text{T,p}} \\ & G_{\text{TOT}} &= N_{\text{W}} \mu_{\text{W}}^{0} + N_{\text{S}} \epsilon_{\text{S}} - T S_{\text{MX}} \end{split}$$

water G + solute energy + mixing entropy contribution Mixing entropy contribution:

- -independent solute molecules (ideal gas)
- -lattice model: $N_w + N_s$ is a total number of lattice sites

$$W(N_w,N_s) = (N_w + N_s)! / (N_w! N_s!)$$

$$S = k_{B} \ln W = \ln[(N_{W} + N_{S})! / (N_{W}! N_{S}!)] \approx$$

$$- k_{B} \{N_{W} \ln[N_{W}/(N_{W} + N_{S})] + N_{S} \ln[N_{S}/(N_{W} + N_{S})]\} =$$

$$- k_{B} [N_{W} \ln(1 - N_{S}/N_{W}) + N_{S} \ln(N_{S}/N_{W})]$$

Taking into account the Taylor expansion of $ln(1+x) \approx x$, we get:

$$\begin{split} \mathbf{S}_{MX} &= - \mathbf{k}_{B} \left[\mathbf{N}_{S} \ln(\mathbf{N}_{S}/\mathbf{N}_{W}) - \mathbf{N}_{S} \right] \\ \mathbf{G}_{TOT} &\left(\mathbf{T}, \mathbf{p}, \mathbf{N}_{W}, \mathbf{N}_{S} \right) = \mathbf{N}_{W} \mu_{W}^{0} + \mathbf{N}_{S} \boldsymbol{\epsilon}_{S} (\mathbf{T}, \mathbf{p}) + \\ & \mathbf{k}_{B} \mathbf{T} (\mathbf{N}_{S} \ln(\mathbf{N}_{S}/\mathbf{N}_{W}) - \mathbf{N}_{S}) \\ \boldsymbol{\mu}_{S} (\mathbf{T}, \mathbf{p}) = \boldsymbol{\epsilon}_{S} (\mathbf{T}, \mathbf{p}) + \mathbf{k}_{B} \mathbf{T} \ln(\mathbf{c}/\mathbf{c}_{0}) \end{split}$$

or in a general form expressed in concentration c=N/V:

$$\mu_{i} = \mu_{i0} + k_{B}T \ln(c_{i}/c_{i0})$$

Osmotic pressure Is an Entropic Effect

- → consider a cell in an aqueous environment exchanging material with a solution (intake of food, excretion of waste)
- \rightarrow chemical potential difference proportional to ΔG :

$$\Delta G = (\mu_1 - \mu_2) dN \le 0$$
 (spontaneous)

- → cell with a crowded environment of biomolecules: tendency of almost all components to move out causes a mechanical pressure called *osmotic pressure*
- → lipid membranes with ion channels to regulate ion concentration
- \rightarrow calculate osmotic pressure due to a dilute solution of $N_{_{\rm S}}$ molecules

Osmotic pressure on a semipermeable membrane, which only allows water molecules through

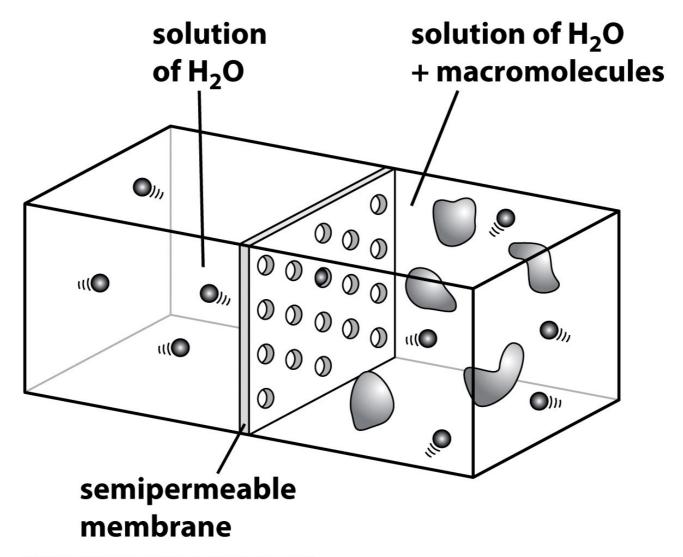


Figure 6.21 Physical Biology of the Cell (© Garland Science 2009)

$$\mu_{W}(T,p) = (\partial G_{TOT} / \partial N_{W})_{T,p} =$$

$$\mu_{W}^{0}(T,p) - k_{B}T N_{S}/N_{W}$$

For both sides of the membrane in equilibrium:

$$\mu_{W}^{0}(T,p_{1}) = \mu_{W}^{0}(T,p_{2}) - k_{B}T N_{S}/N_{W}$$

Expand the chemical potential at \mathbf{p}_2 around the \mathbf{p}_1 value:

$$\mu_W^{\ 0}(T,p_2)\approx \mu_W^{\ 0}(T,p_1)+(\partial\mu_W^{\ 0}/\partial p)\,(p_2-p_1)$$
 and consider that
$$(\partial\mu_W^{\ 0}/\partial p)=v=N_W^{\ V}V \text{ is volume per water molecule so that}$$

$$(\mathbf{p}_2 - \mathbf{p}_1) = \mathbf{k}_B \mathbf{T} \ \mathbf{N}_S / \mathbf{V}$$

Measuring interstrand interactions in DNA using osmotic pressure

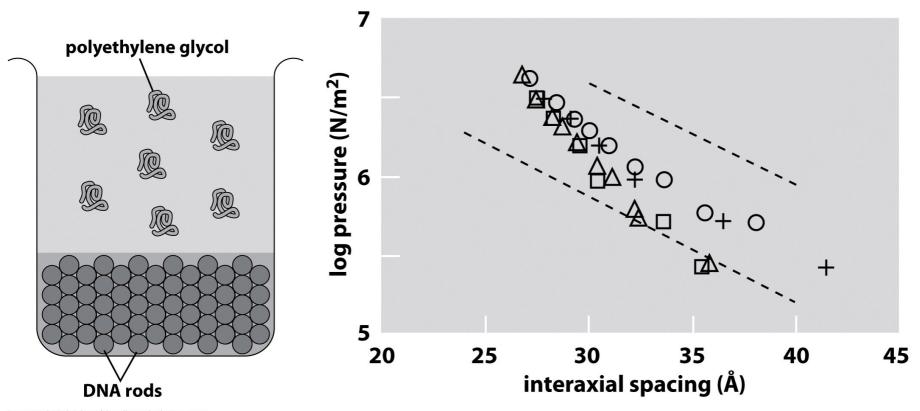


Figure 6.22 Physical Biology of the Cell (© Garland Science 2009)

Figure 6.23 Physical Biology of the Cell (© Garland Science 2009)

pressure:
$$p(d_S) = F_0 \exp(-d_S/c)$$

 $\boldsymbol{d}_{\boldsymbol{S}}$... interstrand spacing; $\boldsymbol{F}_{\boldsymbol{0}}$ depends on the ionic solution

Law of Mass Action and Equilibrium Constants (Chemical Reactions)

→ chemical equilibrium between A, B and their complex AB:

$$A + B \leftrightarrow AB$$

- → final equilibrium independent of whether we start with only A and B, or with a high concentration of AB and no A or B
- $\rightarrow N_A$, N_B , N_{AB} ... number of A, B, and AB molecules
- → In equilibrium: dG = 0 $0 = (\partial G/\partial N_A) dN_A + (\partial G/\partial N_B) dN_B + (\partial G/\partial N_{AB}) dN_{AB}$
- → A more convenient and general expression:

$$\sum_{i \neq i}^{N} \mu_{i} dN_{i} = 0$$

Stoichiometric coefficients for each of the reactants are defined as:

$$v_i = \pm 1$$

depending on whether the number of particles of the i-th type increases or decreases during the reaction:

$$\sum_{i \neq i}^{N} \mu_{i} \nu_{i} = 0$$

$$\sum_{i \neq i}^{N} \mu_{i0} \nu_{i} = -k_{B} T \sum_{i \neq i}^{N} \ln(c_{i}/c_{i0})^{\nu i}$$

or

$$-\beta \sum_{i=1}^{N} \mu_{i0} \nu_{i} = \ln \left[\prod_{i=1}^{N} (c_{i}/c_{i0})^{\nu i} \right]$$

or

$$\prod_{i \neq i} {}^{N} c_{i}^{\nu i} = \left(\prod_{i \neq i} {}^{N} c_{i0}^{\nu i}\right) \exp(-\beta \sum_{i \neq i} {}^{N} \mu_{i0} \nu_{i})$$

where we define the equilibrium constant $\mathbf{K}_{\mathbf{e}\mathbf{n}}$:

$$\mathbf{K}_{eq} = \left(\prod_{i \neq i}^{N} \mathbf{c}_{i0}^{vi}\right) \exp(-\beta \sum_{i \neq i}^{N} \mu_{i0} v_{i})$$

$$K_d = 1/K_{eq}$$
 ... dissociation constant

In our case of the reaction $A + B \leftrightarrow AB$, we can express K_d as:

$$\mathbf{K}_{\mathbf{d}} = \prod_{\mathbf{i} \in \mathbf{d}} \mathbf{C}_{\mathbf{i}}^{\mathbf{N}} \mathbf{C}_{\mathbf{i}}^{-\mathbf{v} \mathbf{i}} = \mathbf{C}_{\mathbf{A}} \mathbf{C}_{\mathbf{B}} / \mathbf{C}_{\mathbf{A}\mathbf{B}}$$

Example: total concentration: 50 µM

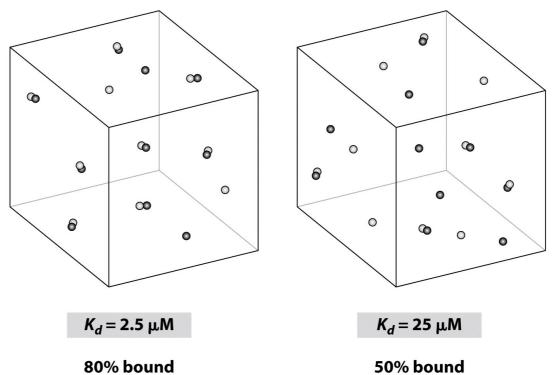


Figure 6.24 Physical Biology of the Cell (© Garland Science 2009)

Application to Ligand-Receptor Binding:

$$L + R \leftrightarrow LR$$

 $K_d = [L][R]/[LR] \text{ or } [LR] = [L] [R]/K_d$

Binding probability:

A natural interpretation of K_d : K_d is the concentration at which the receptor has a probability of ½ of being occupied by a ligand. Based on out prior result, we can express it in terms of lattice

model parameters as:
$$K_d = v^{-1} \exp(\beta \Delta \epsilon)$$

Important: \mathbf{K}_{d} depends on the concentration of *free* ligands not their *total* concentration!

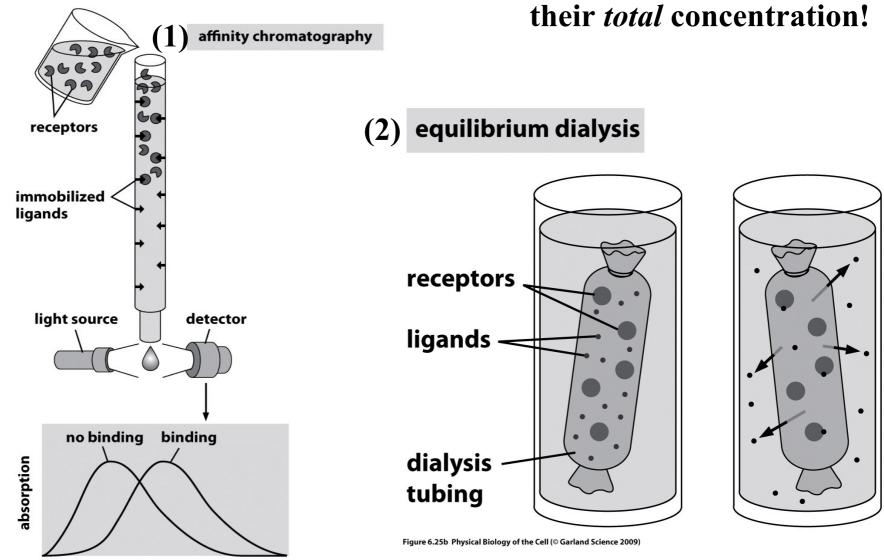
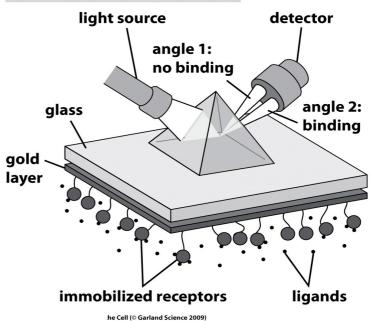


Figure 6.25a Physical Biology of the Cell (© Garland Science 2009)

time

10/11/2011

(C) surface plasmon resonance





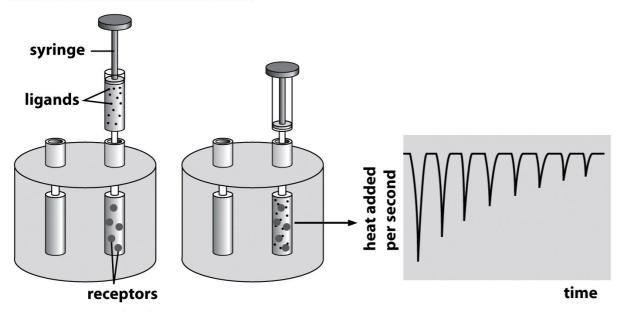


Figure 6.25d Physical Biology of the Cell (© Garland Science 2009) $10/11/2011\,$

Cooperative Ligand-Receptor Binding: The Hill Function

Biological function: an on-off switch behavior triggered by Binding of a ligand to a receptor involves a cooperative (all-or-none) mechanism:

The larger the n, the sharper the binding curve (probability of binding versus ligand concentration)

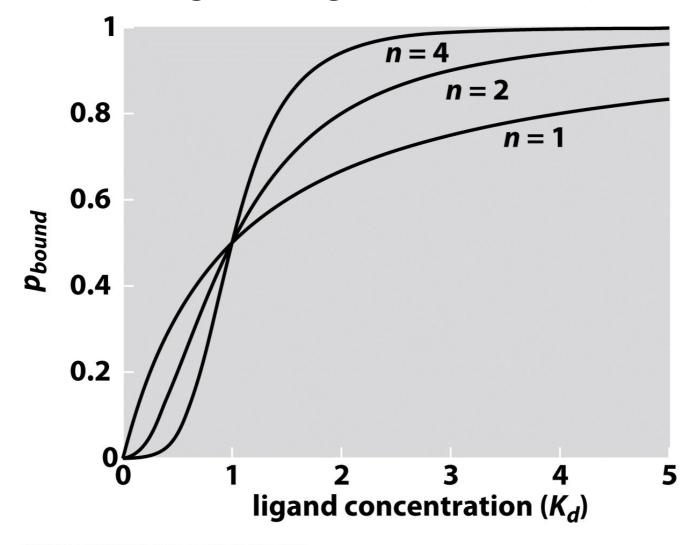


Figure 6.27 Physical Biology of the Cell (© Garland Science 2009)