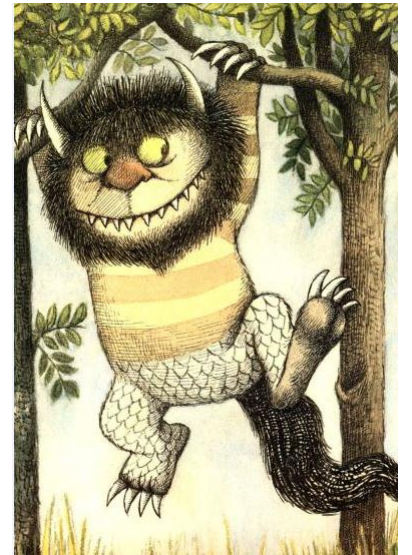
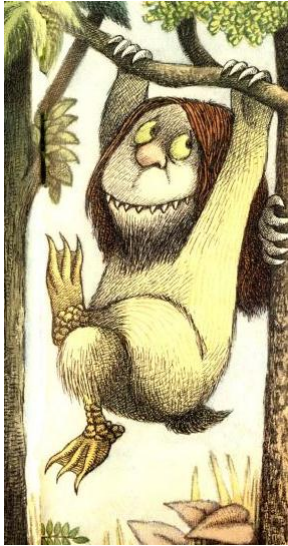


**Damped expectations for
the second generation:
controlling the chaos**









Theorist



Theorist

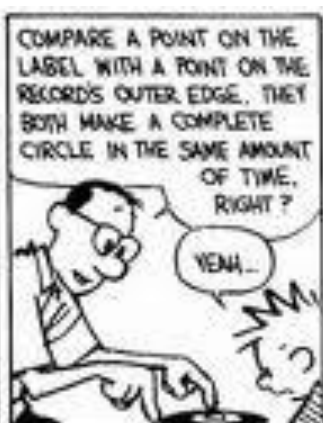


Experimentalist

or



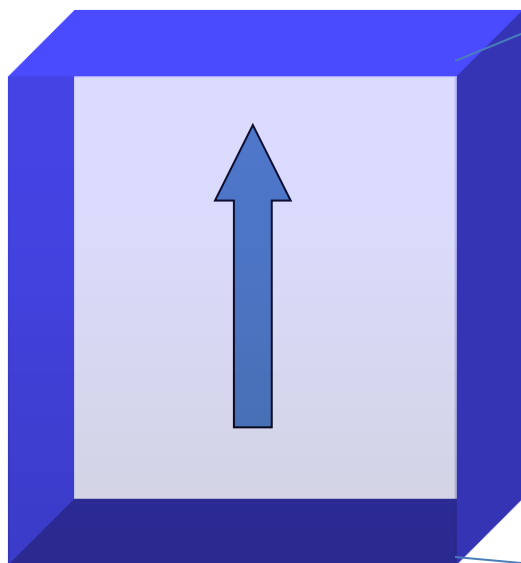




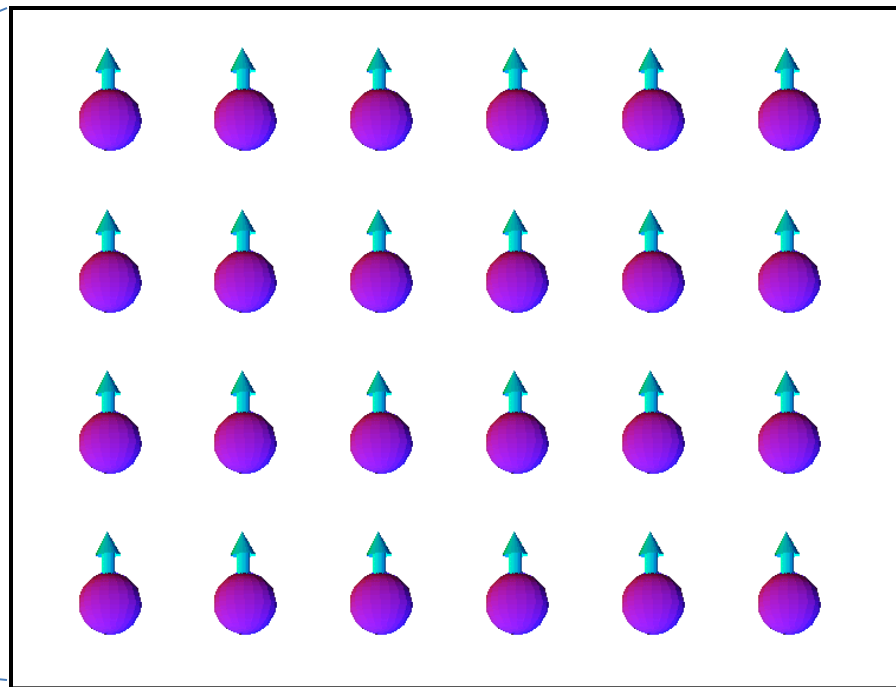




total magnetization



macroscopic picture

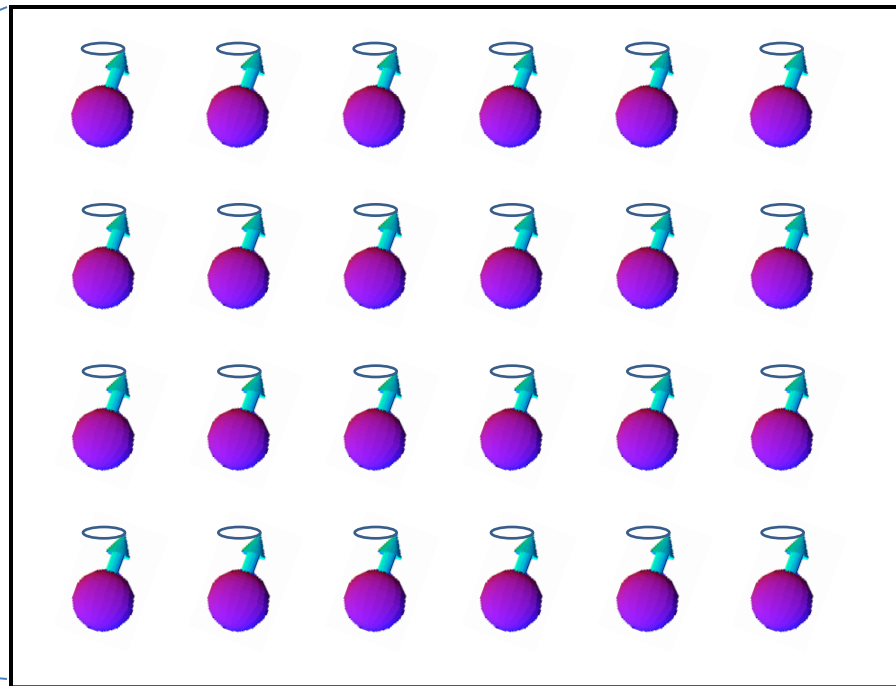


microscopic picture

total magnetization



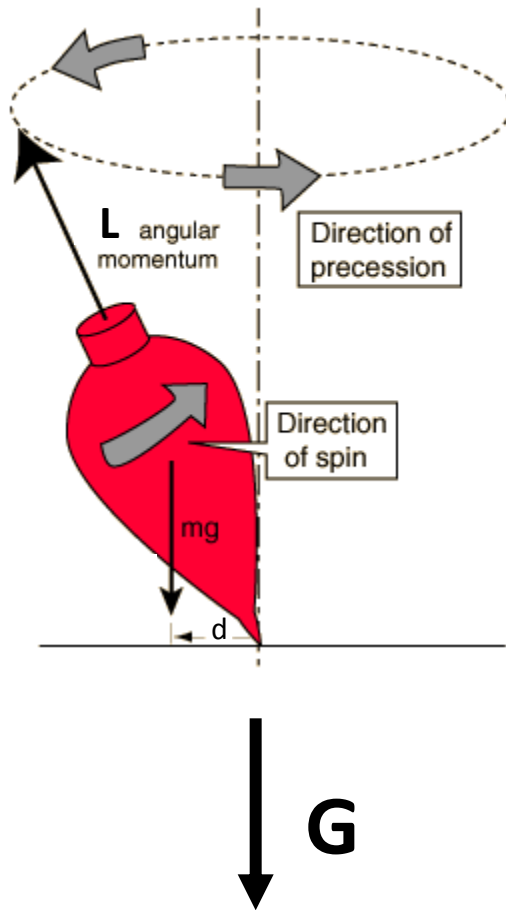
macroscopic picture



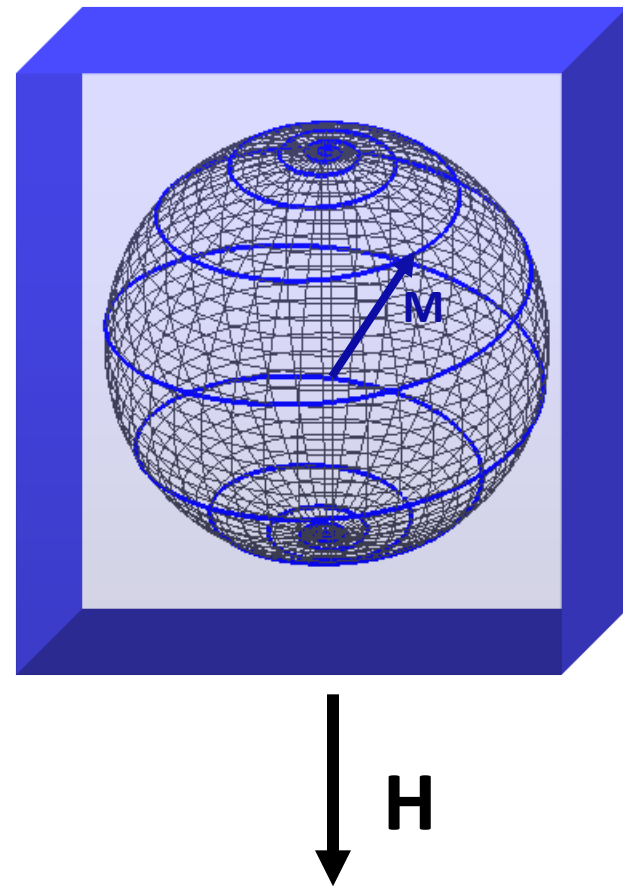
microscopic picture

Field driven magnetodynamics

$$\begin{aligned} d\mathbf{L} / dt &= \mathbf{d} \times m\mathbf{g} && \text{precession} \\ &+ \alpha \hat{\mathbf{L}} \times \dot{\mathbf{L}} && \text{damping} \end{aligned}$$



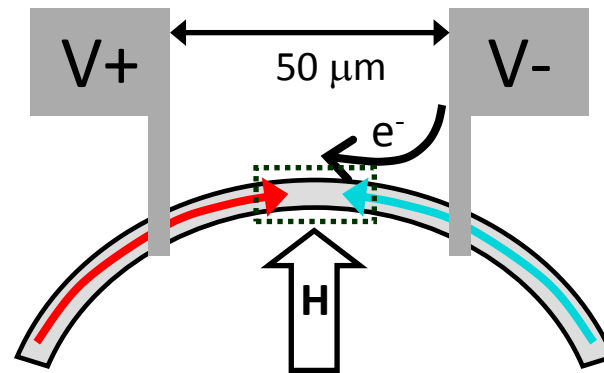
$$\begin{aligned} d\mathbf{M} / dt &= -\mathbf{M} \times \gamma\mathbf{H} && \text{precession} \\ &+ \alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}} && \text{damping} \end{aligned}$$



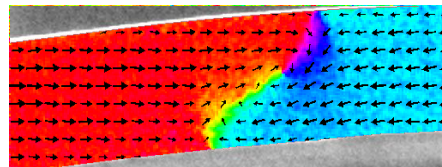
Equation of Motion (damped-driven oscillator)

$$\dot{\mathbf{M}} = \underbrace{-\mathbf{M} \times \gamma \mathbf{H}}_{\text{precession}} + \underbrace{\alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}}_{\text{damping}}$$

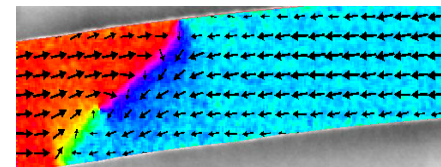
Current-driven magnetodynamics



Before current



After current



Current driven magnetodynamics

Wide domain wall

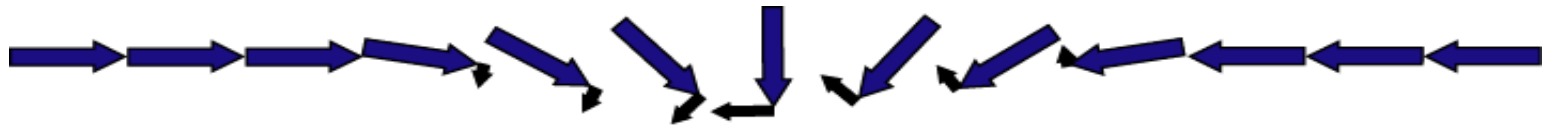
Adiabatic spin-transfer torque



Electron flow



If spins follow magnetization direction



Reaction torque on magnetization



Domain wall translates



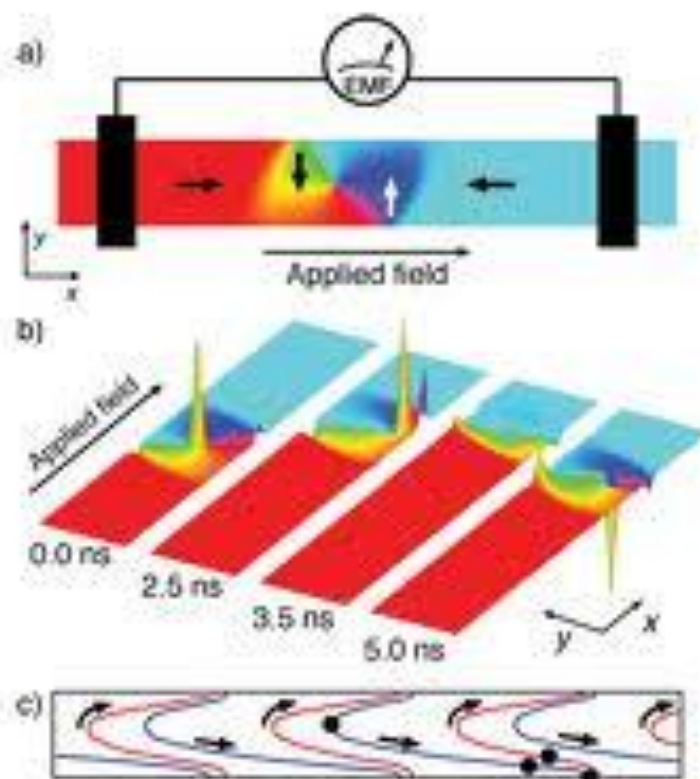
$$\mathbf{v}_s = \frac{-P\mu_B j}{eM_s}$$

Equation of Motion

Field (H) driven terms

Current (v_s) driven terms

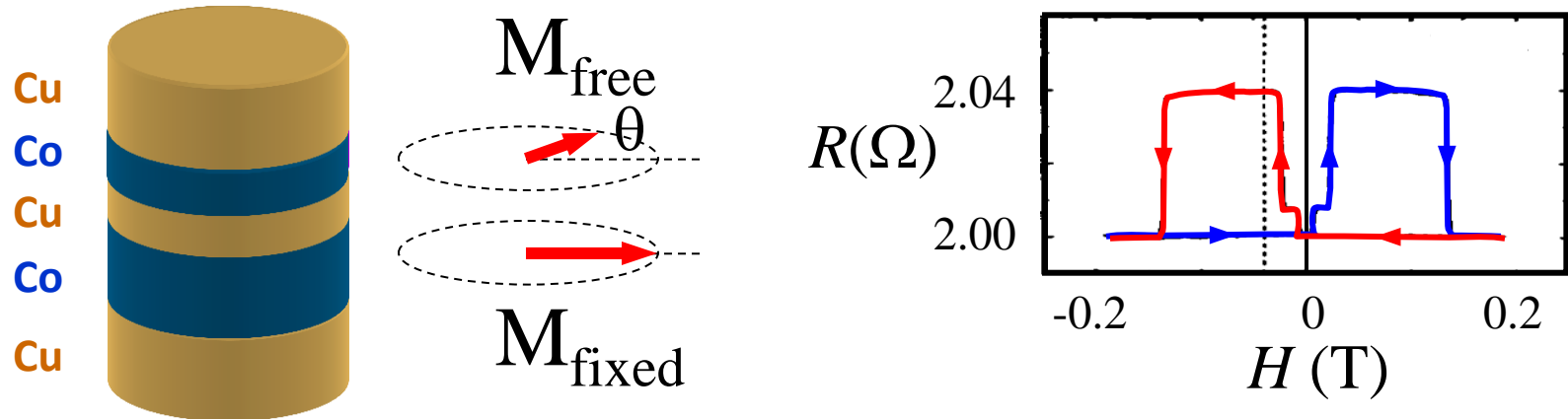
$$\dot{\mathbf{M}} = \underbrace{-\mathbf{M} \times \gamma \mathbf{H}}_{\text{precession}} + \underbrace{\alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}}_{\text{damping}} + \underbrace{-(\mathbf{v}_s \cdot \nabla) \mathbf{M}}_{\text{adiabatic spin torque}} + \underbrace{\beta \hat{\mathbf{M}} \times (\mathbf{v}_s \cdot \nabla) \mathbf{M}}_{\text{non-adiabatic spin torque}}$$





Magnetic state affects current

Giant Magnetoresistance



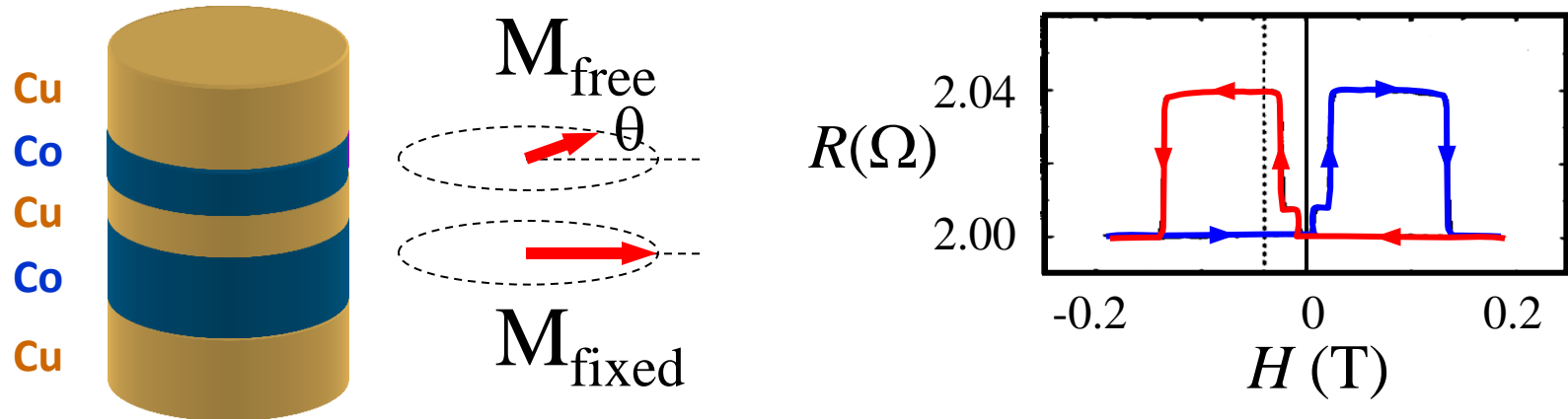
Albert Fert



Peter Grünberg

Magnetic state affects current

Giant Magnetoresistance



Albert Fert



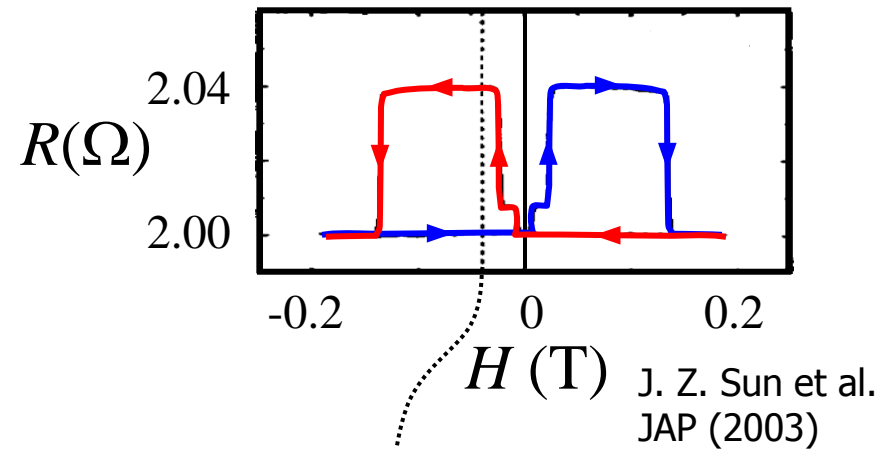
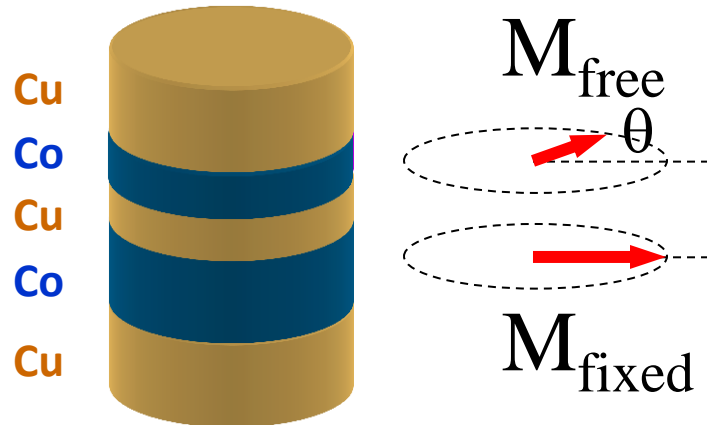
Nobel prize
2007



Peter Grünberg

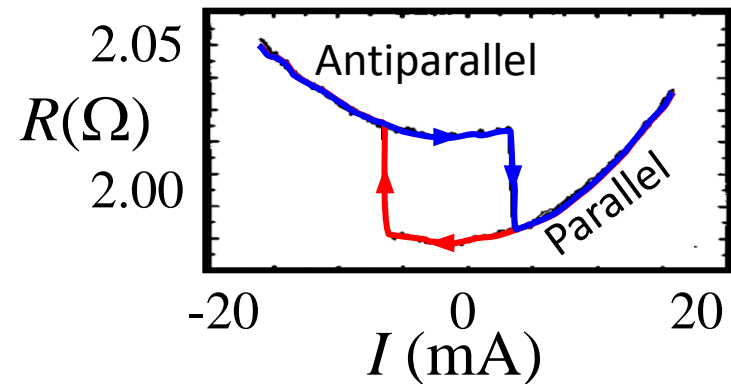
Currents control magnetic state

Giant Magnetoresistance



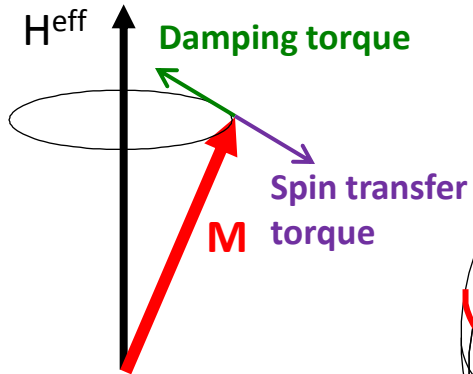
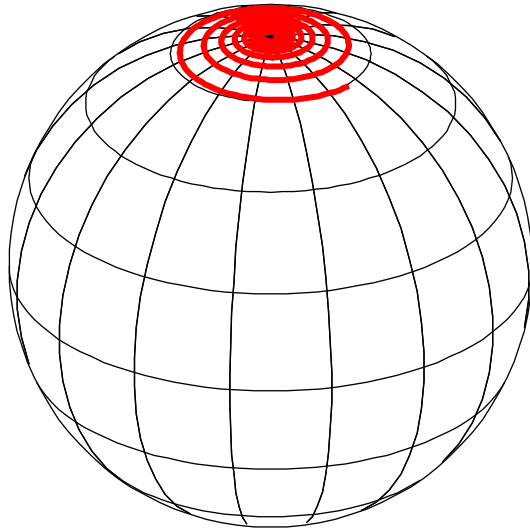
Spin-Transfer Torque

Spin currents can switch magnetization!

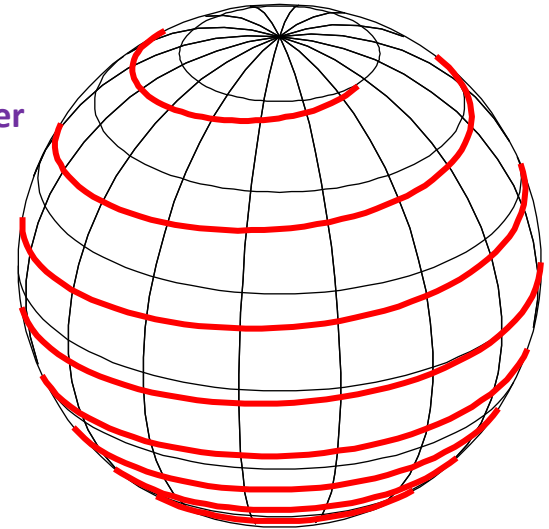


Current-induced switching dynamics

Low current
→ damped motion

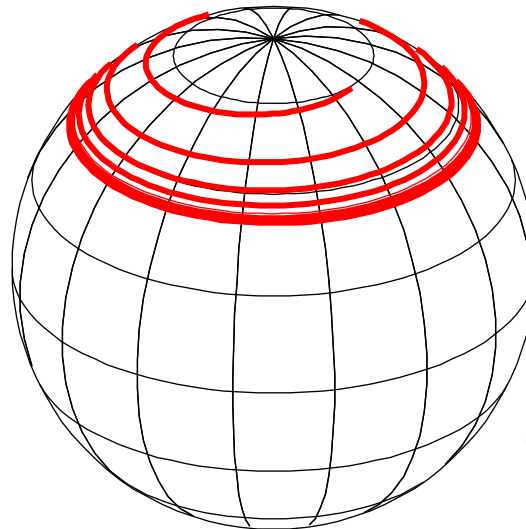


High current
→ switching



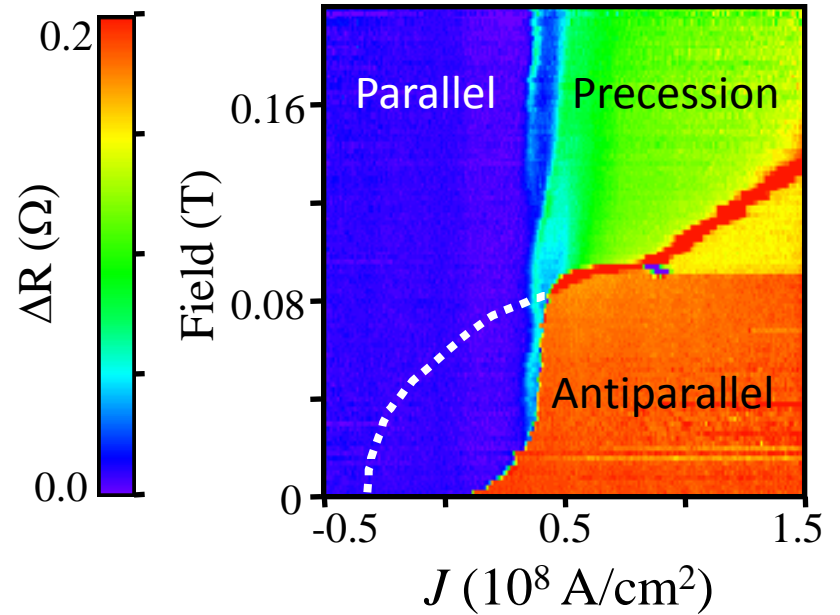
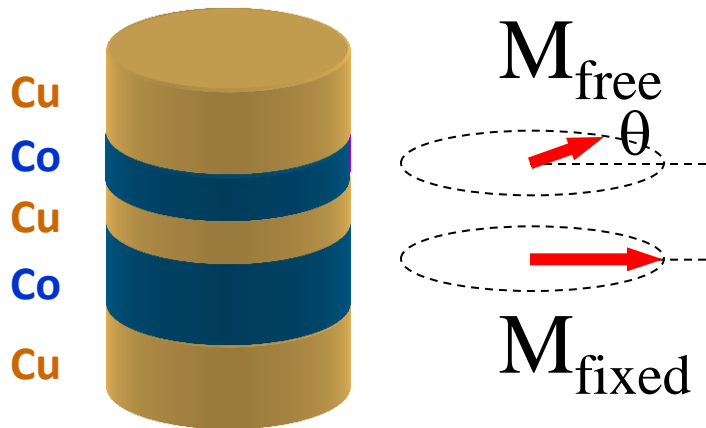
Better MRAM

High current, high field
→ stable precession



**Current-tunable
microwave oscillators**

Complex phase diagram



Equation of Motion

Field (H) driven terms

Current (v_s) driven terms

$$\dot{\mathbf{M}} = \underbrace{-\mathbf{M} \times \gamma \mathbf{H}}_{\text{precession}} + \underbrace{\alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}}_{\text{damping}} + \underbrace{-(\mathbf{v}_s \cdot \nabla) \mathbf{M}}_{\text{adiabatic spin torque}} + \underbrace{\beta \hat{\mathbf{M}} \times (\mathbf{v}_s \cdot \nabla) \mathbf{M}}_{\text{non-adiabatic spin torque}}$$

well understood, conservative

less well understood, dissipative

well understood, conservative

less well understood, dissipative

Equation of Motion

We want to understand the origin of these parameters

$$\dot{\mathbf{M}} = \underbrace{-\mathbf{M} \times \gamma \mathbf{H}}_{\text{precession}} + \underbrace{\alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}}_{\text{damping}} + \underbrace{-(\mathbf{v}_s \cdot \nabla) \mathbf{M}}_{\text{adiabatic spin torque}} + \underbrace{\beta \hat{\mathbf{M}} \times (\mathbf{v}_s \cdot \nabla) \mathbf{M}}_{\text{non-adiabatic spin torque}}$$

The diagram illustrates the equation of motion for magnetization dynamics, $\dot{\mathbf{M}}$, divided into two parts:

- Left part (green dashed box):** Contains the precession term $-\mathbf{M} \times \gamma \mathbf{H}$ and the damping term $\alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}$. The precession term is labeled "well understood, conservative" and "precession". The damping term is labeled "less well understood, dissipative" and "damping".
- Right part (purple dashed box):** Contains the adiabatic spin torque term $-(\mathbf{v}_s \cdot \nabla) \mathbf{M}$ and the non-adiabatic spin torque term $\beta \hat{\mathbf{M}} \times (\mathbf{v}_s \cdot \nabla) \mathbf{M}$. The adiabatic spin torque term is labeled "well understood, conservative" and "adiabatic spin torque". The non-adiabatic spin torque term is labeled "less well understood, dissipative" and "non-adiabatic spin torque".

Equation of Motion

$$\dot{\mathbf{M}} = -\mathbf{M} \times \gamma \mathbf{H} + \alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}$$

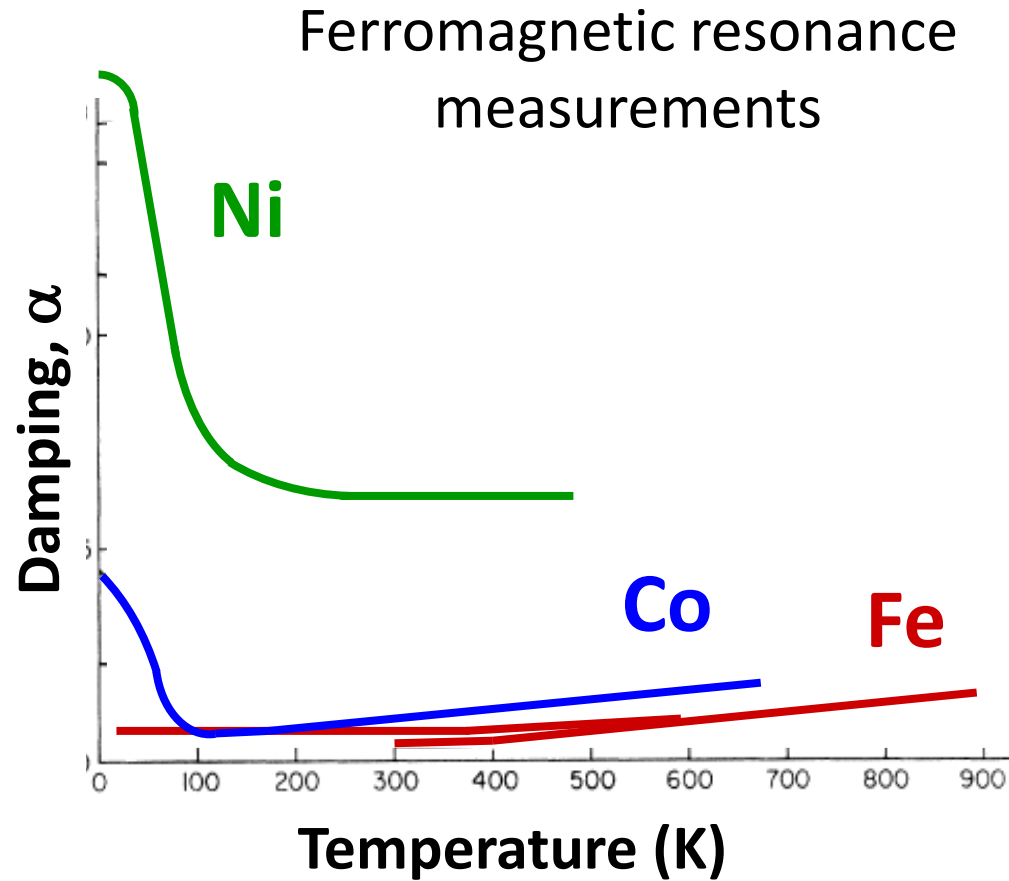
well understood,
conservative

precession

damping

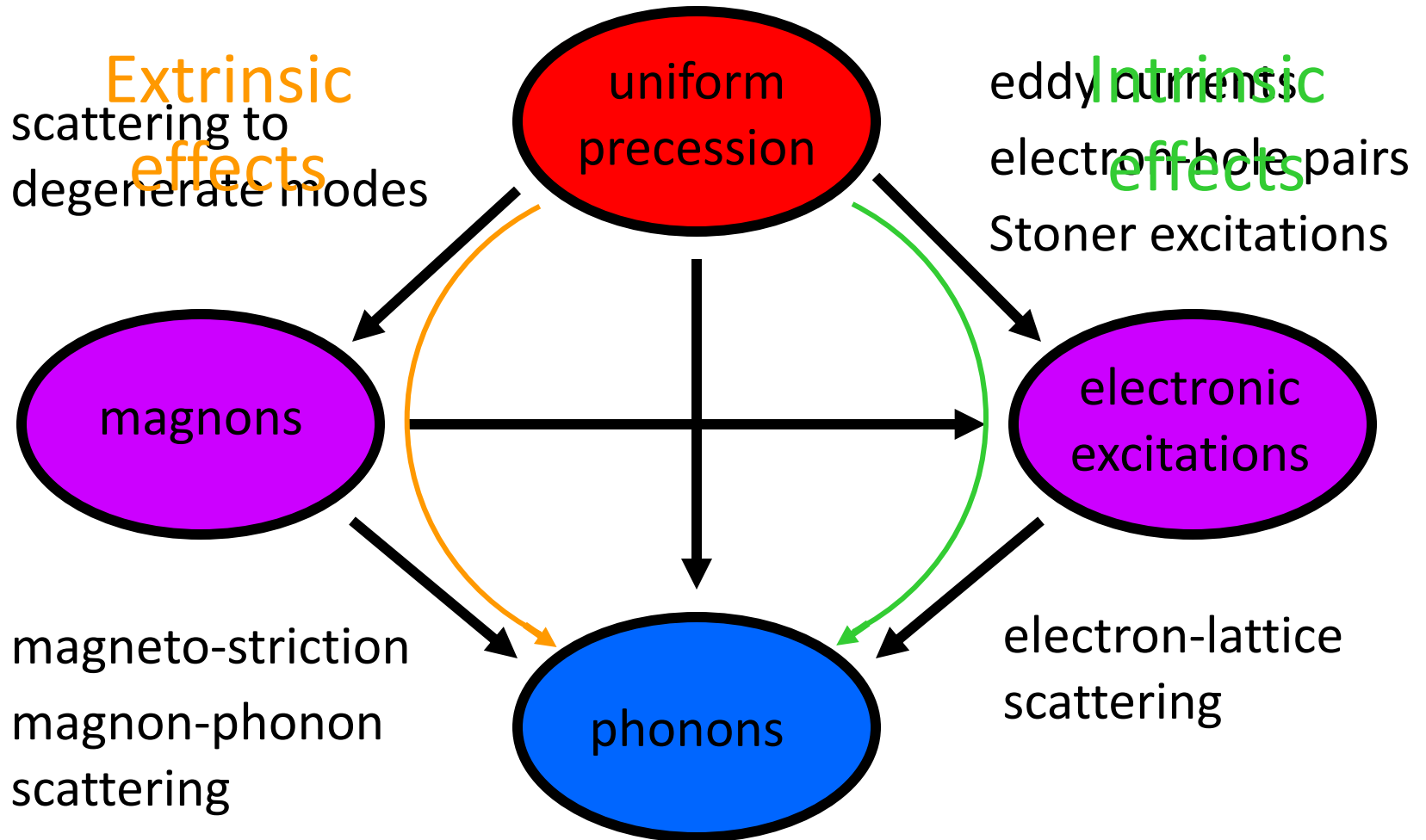
less well understood,
dissipative

Motivating data

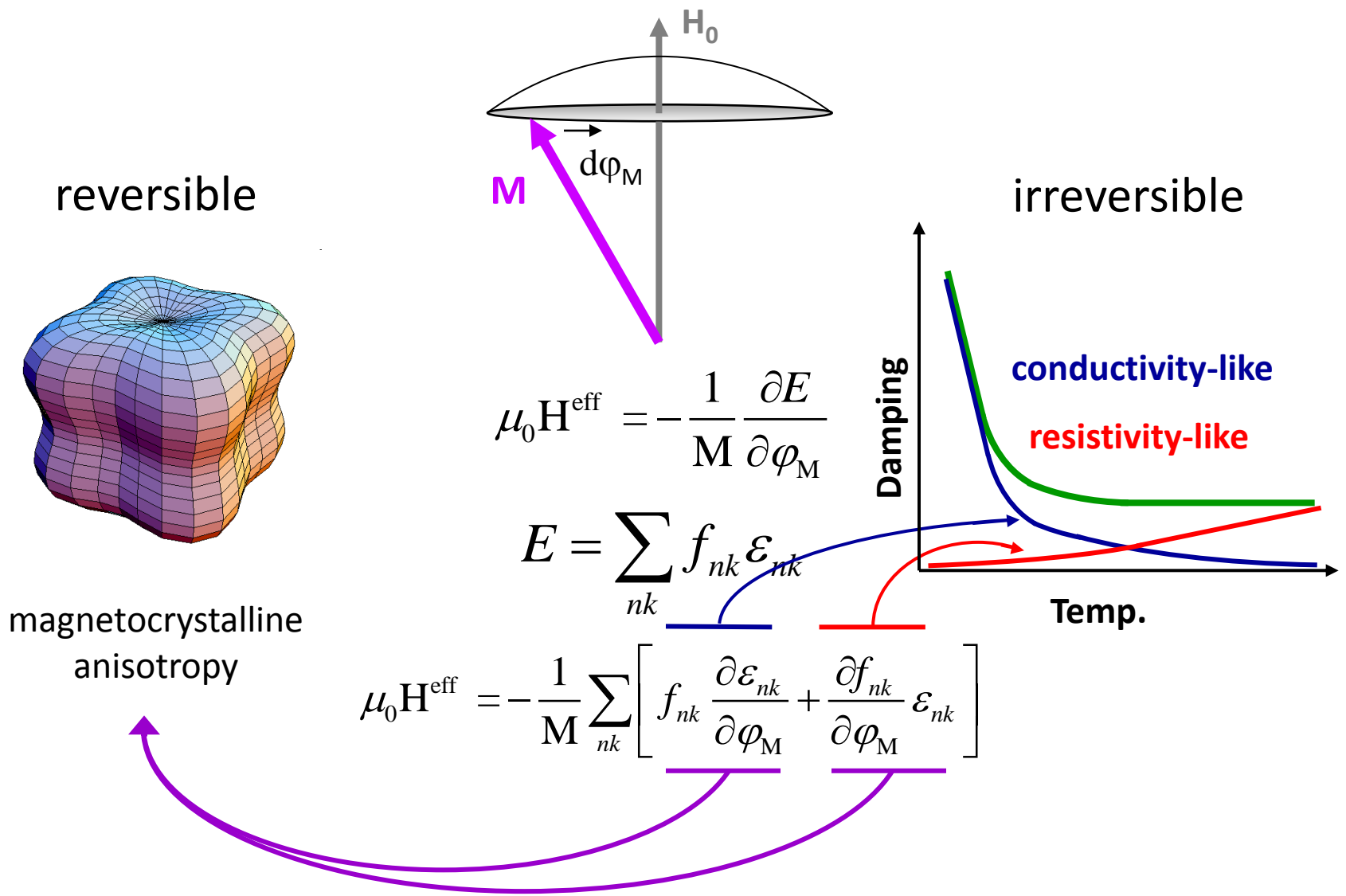


$$\dot{\mathbf{M}} = -\mathbf{M} \times \gamma \mathbf{H} + \alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}$$

Coupling to the environment



Effective Field Model



Fermi Surfaces

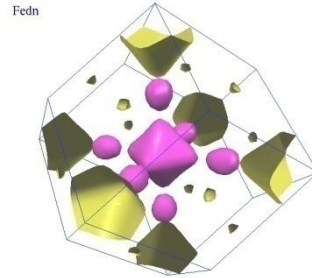
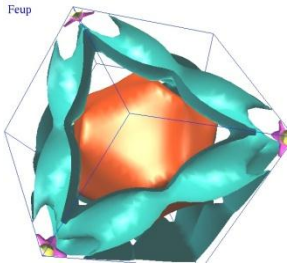
Real Fermi Surfaces

Representative Fermi Surfaces

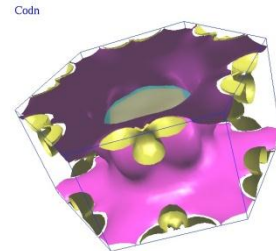
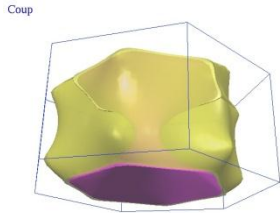
Up Spin

Down Spin

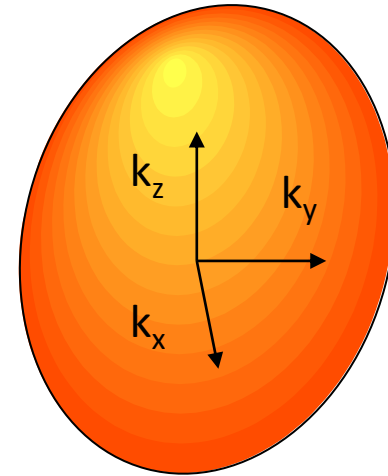
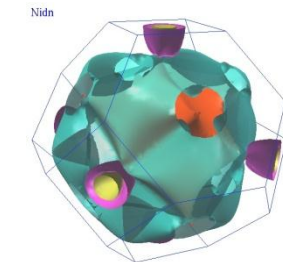
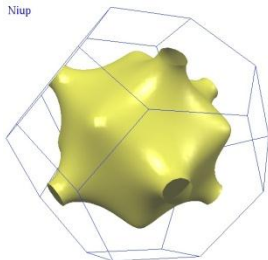
Iron



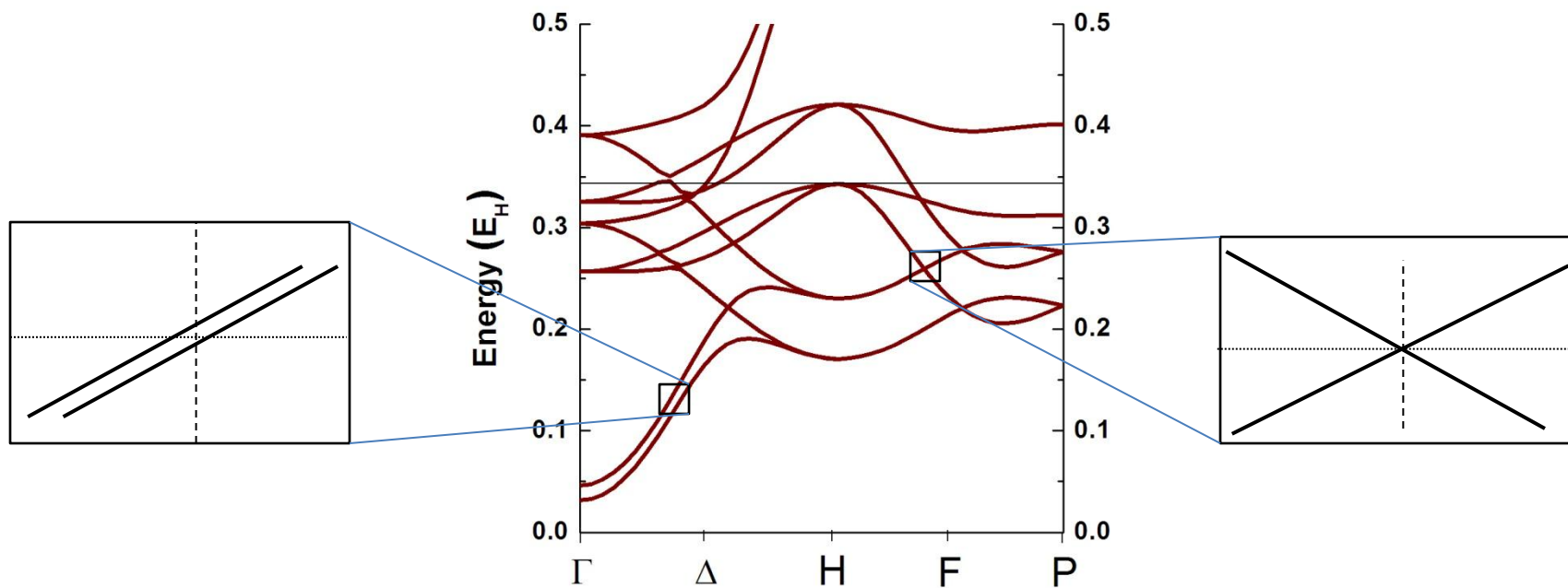
Cobalt



Nickel



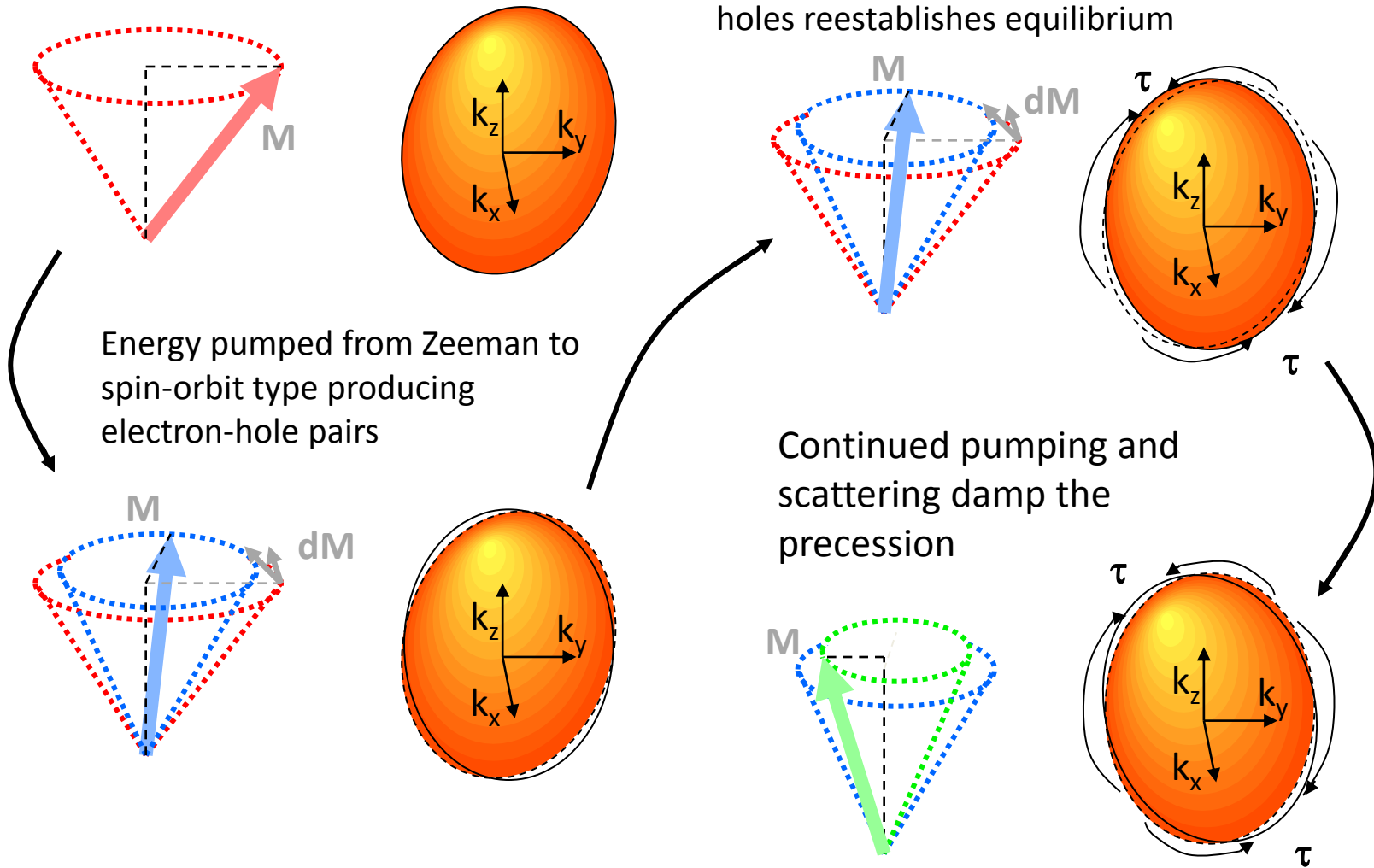
Band Structure



$$\mathbf{H} \varphi_{nk} = \varepsilon_{nk} \varphi_{nk}$$

Breathing Fermi Surface : $H_{so}(t) = -\xi \ell \cdot s(t)$

Scattering of excited electrons into holes reestablishes equilibrium

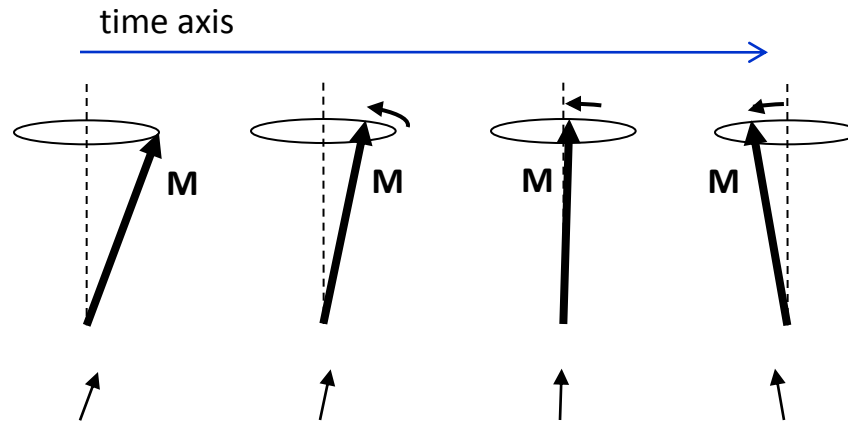


State energies change with precession

Spin-orbit energy depends on spin direction (σ)

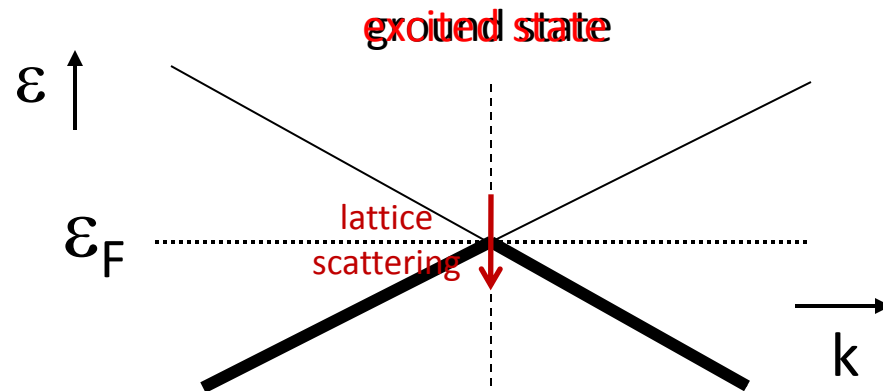
$$H_{so} \approx -\xi \sum_{nk} \ell_{nk} \cdot \sigma_{nk} (\hat{m})$$

Representative spin:

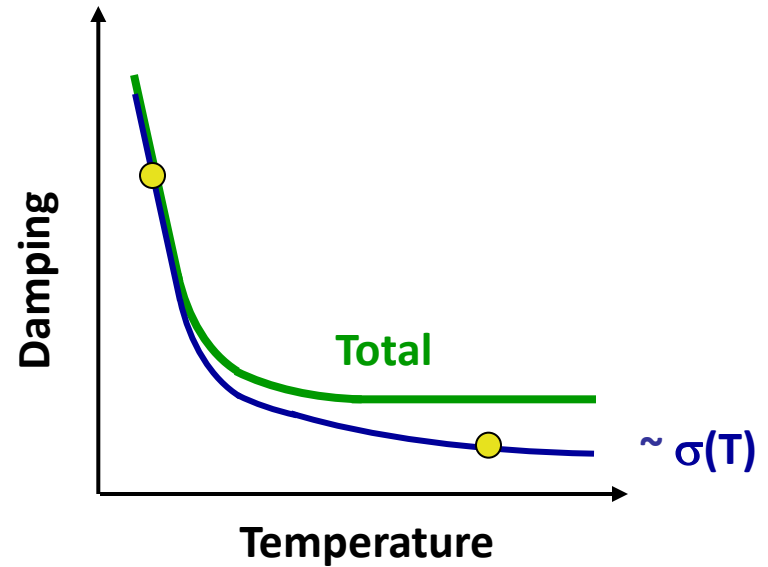
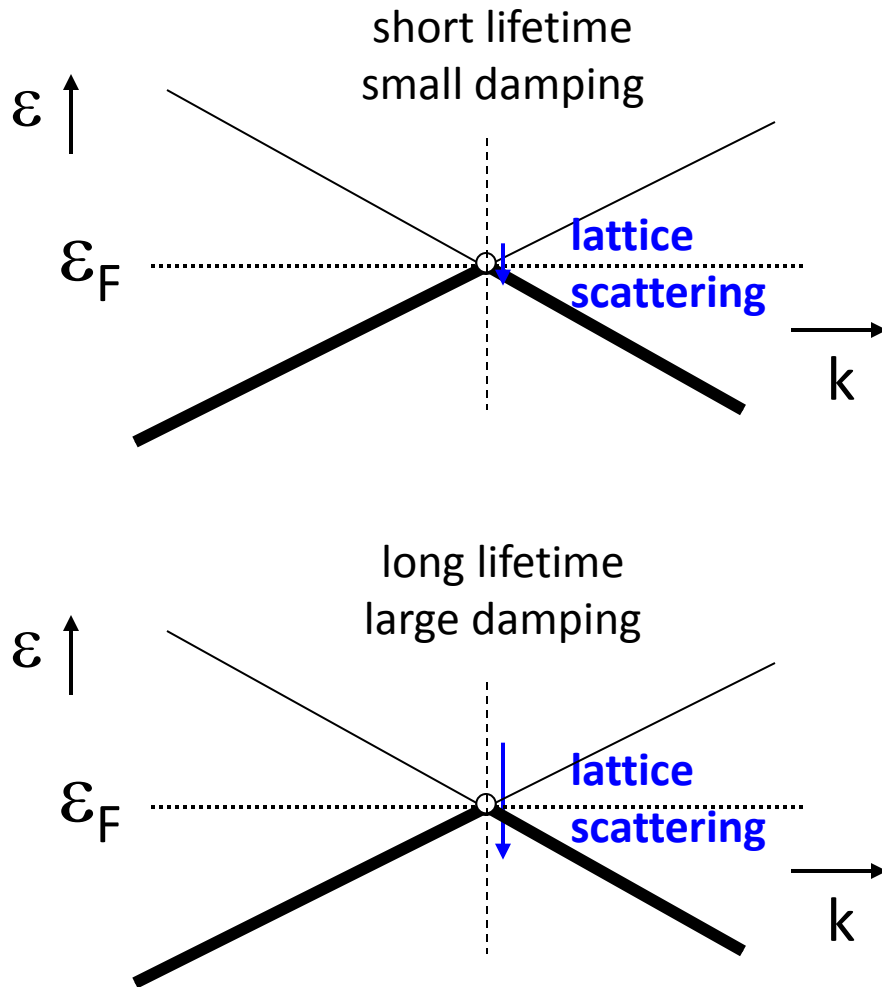


Generation of electron-hole pairs due to energy changes – **breathing Fermi surface**

$$\mu_0 \mathbf{H}^{\text{eff}} = -\frac{1}{M} \sum_{nk} f_{nk} \frac{\partial \varepsilon_{nk}}{\partial \phi_M}$$



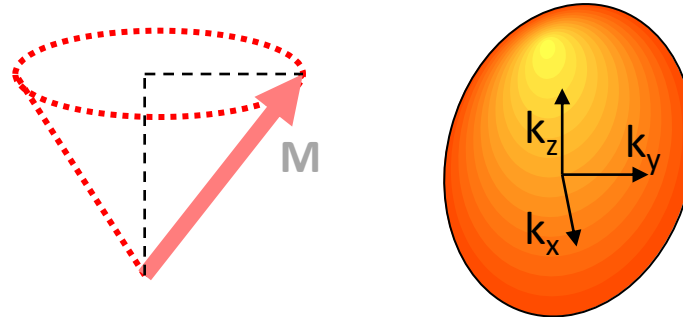
Conductivity-like term *decreases* as scattering rate increases



The **breathing** Fermi surface contribution

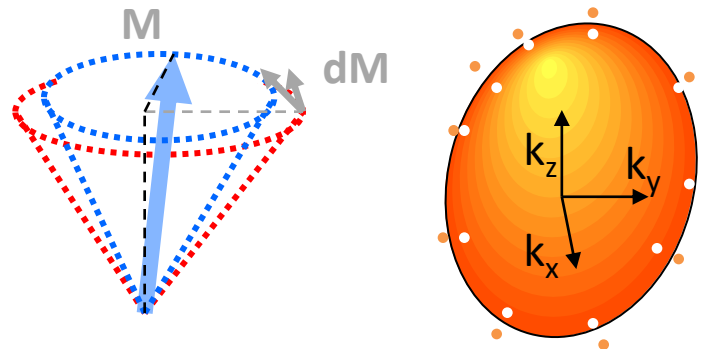
Bubbling Fermi Surface : $V(t) = H_{so}(t) - H_{so}$

$$V(0) = 0$$



The effective perturbation generates electronic excitations over the Fermi surface, reducing the Zeeman energy.

$$V(t) = H_{so}(t) - H_{so}(0)$$



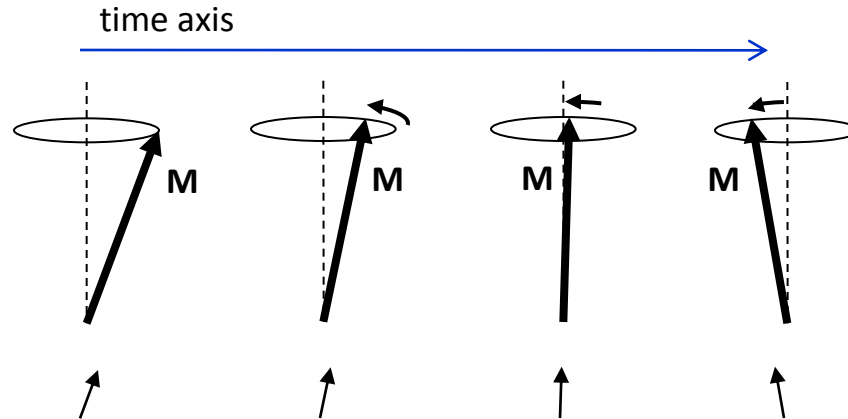
Electron-electron and electron-lattice scattering restores equilibrium.

State populations change with precession

Eigenstates depend on spin direction (σ)

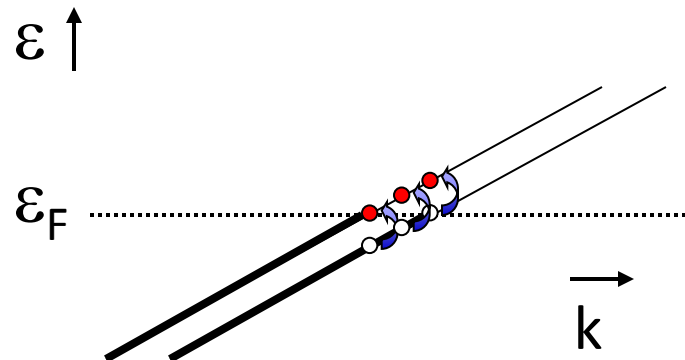
$$V_{so} = H_{so}(\hat{m}) - H_{so}(\hat{z})$$

Representative spin:

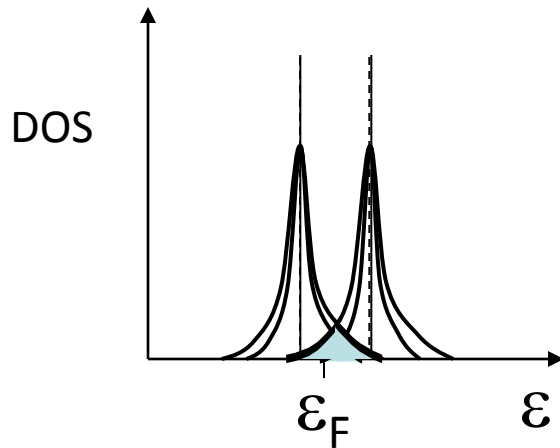
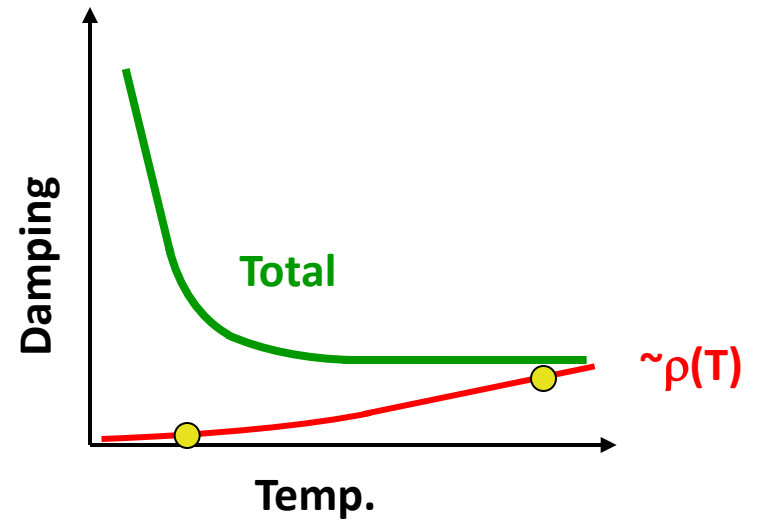
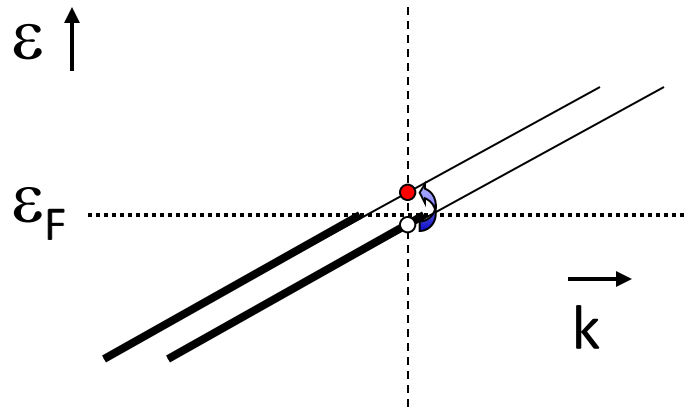


Generation of electron-hole pairs due to population changes – **bubbling Fermi surface**

$$\mu_0 \mathbf{H}^{\text{eff}} = -\frac{1}{M} \sum_{nk} \frac{\partial f_{nk}}{\partial \varphi_M} \varepsilon_{nk}$$



Resistivity-like term *increases* as scattering rate increases



The **bubbling** Fermi surface contribution

Qualitative prediction

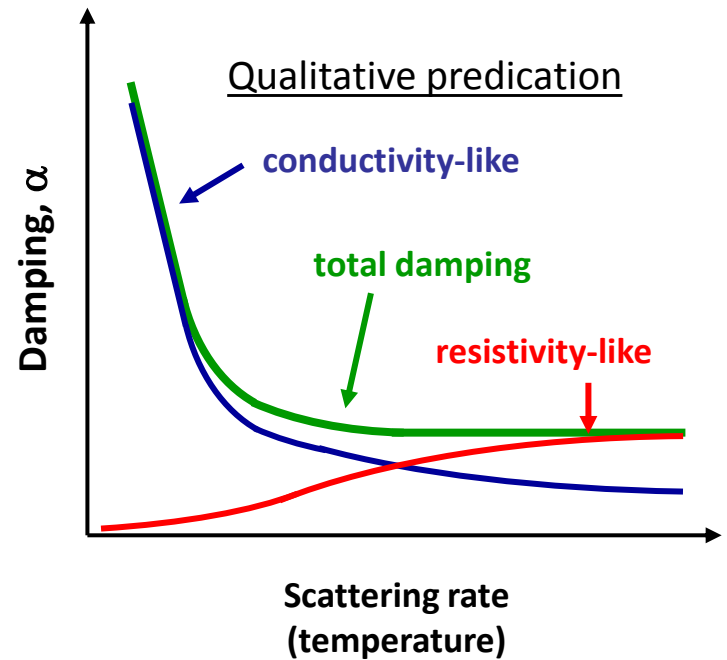
$$\mu_0 \mathbf{H}^{\text{eff}} = -\frac{1}{M_s} \sum_{nk} \left[\underbrace{f_{nk}}_{\text{blue}} \frac{\partial \varepsilon_{nk}}{\partial \varphi_M} + \underbrace{\frac{\partial f_{nk}}{\partial \varphi_M}}_{\text{red}} \varepsilon_{nk} \right]$$

conductivity-like 'breathing' terms

$$\mu_0 \mathbf{H}_{(1)}^{\text{eff}} = -\frac{1}{M_s} \sum_{nk} f_{nk} \frac{\partial \varepsilon_{nk}}{\partial \varphi_M}$$

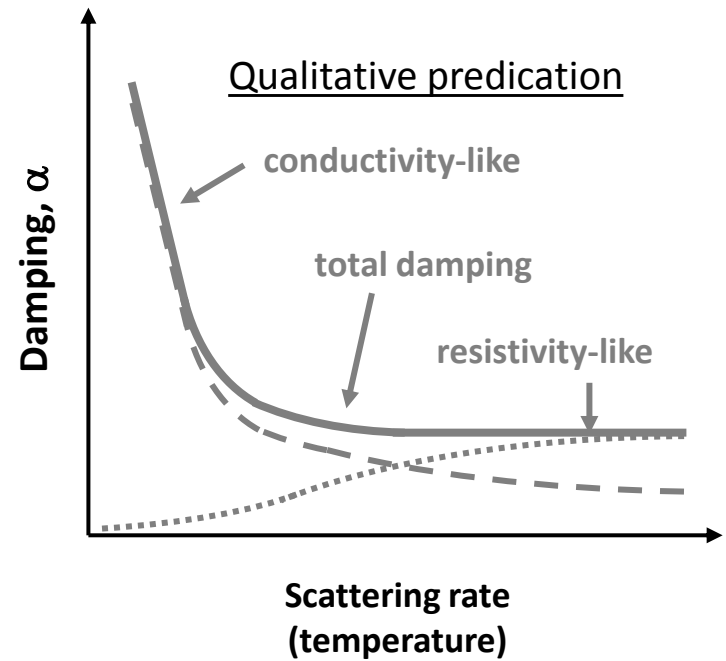
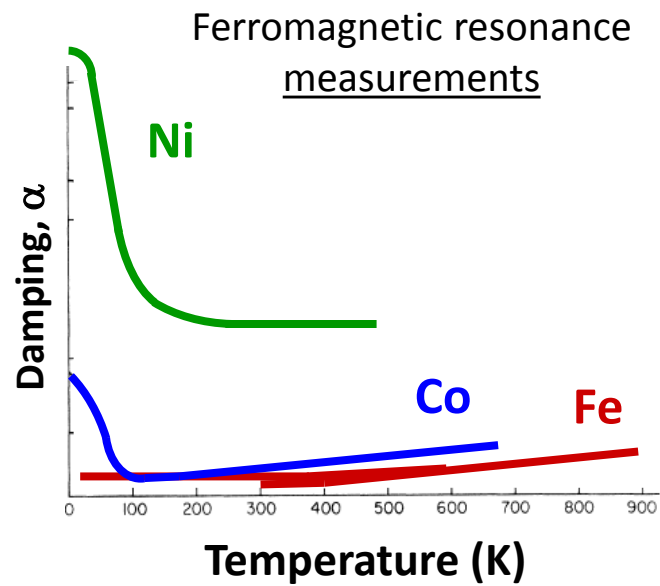
Resistivity-like 'bubbling' terms

$$\mu_0 \mathbf{H}_{(2)}^{\text{eff}} = -\frac{1}{M_s} \sum_{nk} \frac{\partial f_{nk}}{\partial \varphi_M} \varepsilon_{nk}$$



$$\dot{\mathbf{M}} = -\mathbf{M} \times \gamma \mathbf{H} + \alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}$$

Qualitative match with experiment



$$\dot{\mathbf{M}} = -\mathbf{M} \times \gamma \mathbf{H} + \alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}$$

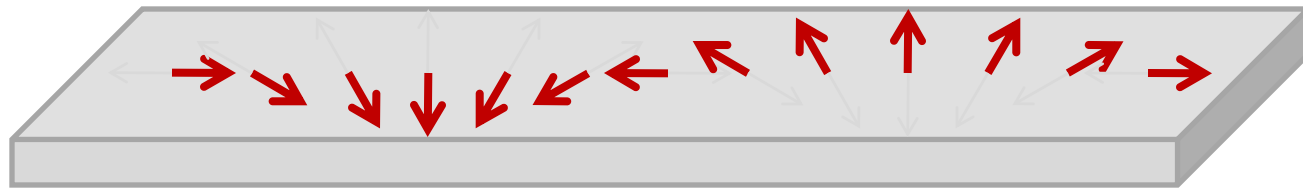
Equation of Motion

$$\dot{\mathbf{M}} =$$

Non-adiabatic STT as extension of damping

$$\alpha \propto \text{Im } \chi(\omega, \mathbf{q}, \mathbf{E} = \mathbf{0})$$

$$\beta \propto \text{Im } \chi(\omega = 0, \mathbf{q}, \mathbf{E})$$



precession

adiabatic spin torque

Conservative
torques

$$\dot{\mathbf{M}} = -\mathbf{M} \times \gamma \mathbf{H} + \underbrace{\alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}}_{\text{damping}} - (\mathbf{v}_s \cdot \nabla) \mathbf{M} + \beta \hat{\mathbf{M}} \times (\mathbf{v}_s \cdot \nabla) \mathbf{M}$$

Dissipative
torques

damping

non-adiabatic spin torque

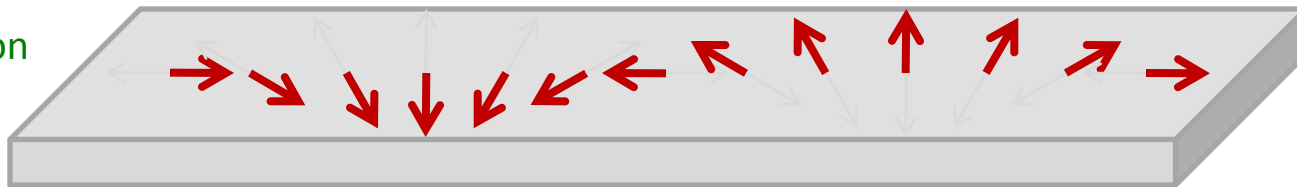
Non-adiabatic STT as extension of damping

$$\alpha \propto \text{Im } \chi(\omega, \mathbf{q}, \mathbf{E} = \mathbf{0})$$

$$\beta \propto \text{Im } \chi(\omega = 0, \mathbf{q}, \mathbf{E})$$

Rotation in time: \rightarrow

precession on
current off



precession

adiabatic spin torque

Conservative
torques

$$\dot{\mathbf{M}} = -\mathbf{M} \times \gamma \mathbf{H} + \underbrace{\alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}}_{\text{damping}} - (\mathbf{v}_s \cdot \nabla) \mathbf{M} + \underbrace{\beta \hat{\mathbf{M}} \times (\mathbf{v}_s \cdot \nabla) \mathbf{M}}_{\text{non-adiabatic spin torque}}$$

Dissipative
torques \rightarrow

damping

non-adiabatic spin torque

Non-adiabatic STT as extension of damping

$$\alpha \propto \text{Im } \chi(\omega, \mathbf{q}, \mathbf{E} = \mathbf{0})$$

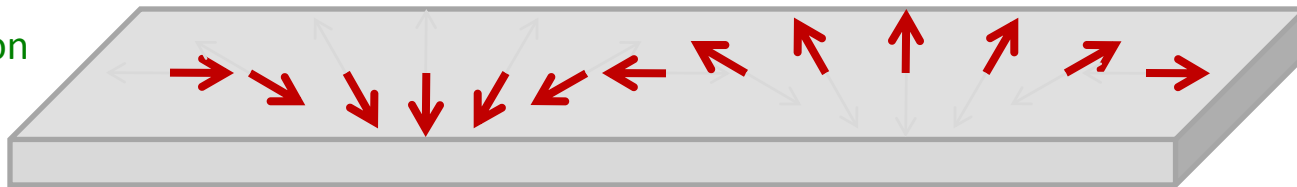
$$\beta \propto \text{Im } \chi(\omega = 0, \mathbf{q}, \mathbf{E})$$

α : all Fermi level electrons contribute

$$\alpha = \sum_{nk} \alpha_{nk}$$

Rotation in time: \rightarrow

precession on
current off



precession

adiabatic spin torque

Conservative
torques

$$\dot{\mathbf{M}} = -\mathbf{M} \times \gamma \mathbf{H} + \underbrace{\alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}}_{\text{damping}} - (\mathbf{v}_s \cdot \nabla) \mathbf{M} + \beta \hat{\mathbf{M}} \times (\mathbf{v}_s \cdot \nabla) \mathbf{M}$$

Dissipative
torques

damping

non-adiabatic spin torque

Non-adiabatic STT as extension of damping

$$\alpha \propto \text{Im } \chi(\omega, \mathbf{q}, \mathbf{E} = \mathbf{0})$$

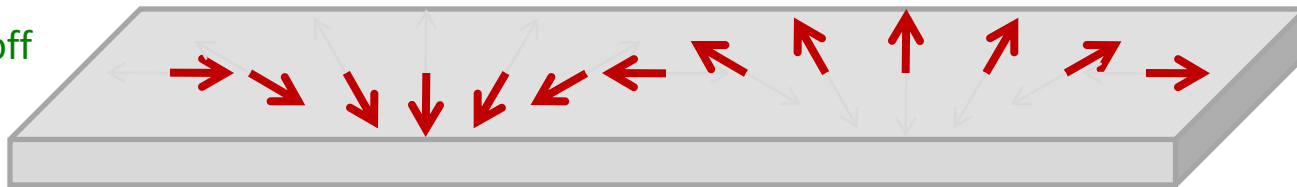
$$\beta \propto \text{Im } \chi(\omega = 0, \mathbf{q}, \mathbf{E})$$

α : all Fermi level electrons contribute

$$\alpha = \sum_{nk} \alpha_{nk}$$

Rotation in space: \rightarrow

precession off
current on



precession

adiabatic spin torque

Conservative
torques

$$\dot{\mathbf{M}} = -\mathbf{M} \times \gamma \mathbf{H} + \underbrace{\alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}}_{\text{damping}} - (\mathbf{v}_s \cdot \nabla) \mathbf{M} + \beta \hat{\mathbf{M}} \times (\mathbf{v}_s \cdot \nabla) \mathbf{M}$$

Dissipative
torques

damping

non-adiabatic spin torque

Non-adiabatic STT as extension of damping

$$\alpha \propto \text{Im } \chi(\omega, \mathbf{q}, \mathbf{E} = \mathbf{0})$$

$$\beta \propto \text{Im } \chi(\omega = 0, \mathbf{q}, \mathbf{E})$$

α : all Fermi level electrons contribute

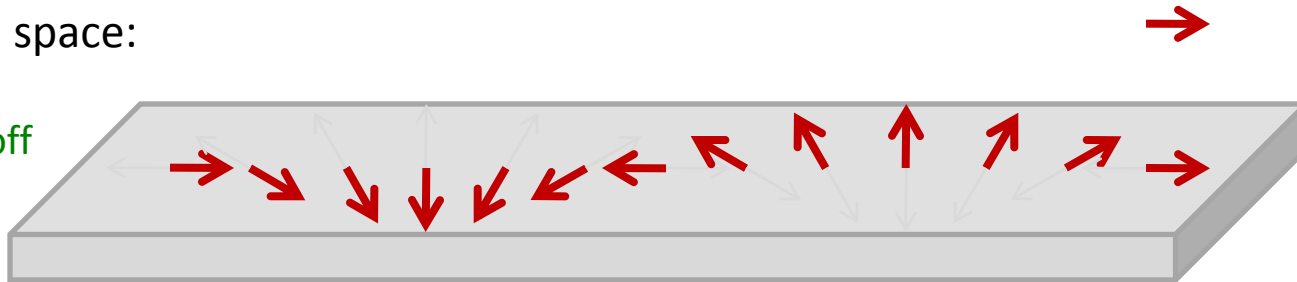
β : Fermi level electron contributions are weighted by their velocity

$$\alpha = \sum_{nk} \alpha_{nk}$$

$$\beta = \sum_{nk} \alpha_{nk} \frac{\mathbf{v}_{nk}}{v_s}$$

Rotation in space:

precession off
current on



precession

adiabatic spin torque

Conservative
torques

$$\dot{\mathbf{M}} = -\mathbf{M} \times \gamma \mathbf{H} + \underbrace{\alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}}_{\text{damping}} - (\mathbf{v}_s \cdot \nabla) \mathbf{M} + \underbrace{\beta \hat{\mathbf{M}} \times (\mathbf{v}_s \cdot \nabla) \mathbf{M}}_{\text{non-adiabatic spin torque}}$$

Dissipative
torques

damping

non-adiabatic spin torque

Non-adiabatic STT as extension of damping

$$\alpha \propto \text{Im } \chi(\omega, \mathbf{q}, \mathbf{E} = \mathbf{0})$$

$$\beta \propto \text{Im } \chi(\omega = 0, \mathbf{q}, \mathbf{E})$$

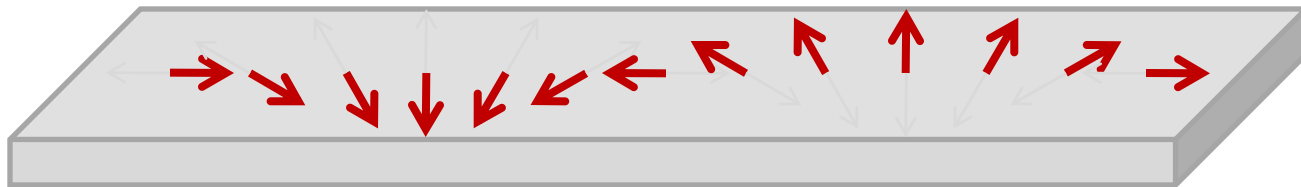
α : all Fermi level electrons contribute

β : Fermi level electron contributions are weighted by their velocity

$$\alpha = \sum_{nk} \alpha_{nk}$$

$$\beta = \sum_{nk} \alpha_{nk} \frac{\mathbf{v}_{nk}}{v_s}$$

$$\frac{d\hat{s}}{dt} = \frac{\partial \hat{s}}{\partial t} + \dot{\mathbf{r}} \cdot \frac{\partial \hat{s}}{\partial \mathbf{r}}$$



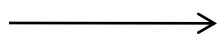
precession

adiabatic spin torque

Conservative torques

$$\dot{\mathbf{M}} = -\mathbf{M} \times \gamma \mathbf{H} + \underbrace{\alpha \hat{\mathbf{M}} \times \dot{\mathbf{M}}}_{\text{damping}} - (\mathbf{v}_s \cdot \nabla) \mathbf{M} + \underbrace{\beta \hat{\mathbf{M}} \times (\mathbf{v}_s \cdot \nabla) \mathbf{M}}_{\text{non-adiabatic spin torque}}$$

Dissipative torques



damping

non-adiabatic spin torque

Qualitative Summary

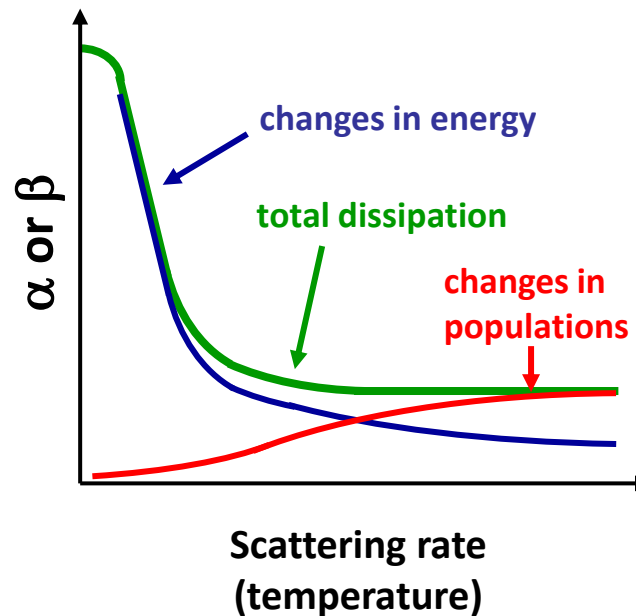
Dissipation occurs through the spin-orbit interaction and electron-lattice scattering

α : precession dissipation parameter

β : spin-transfer torque dissipation parameter

conductivity-like
breathing terms

resistivity-like
bubbling terms



$$\mu_0 \mathbf{H}^{\text{eff}} = -\frac{1}{M} \frac{\partial}{\partial \varphi_M} \sum_{nk} f_{nk} \varepsilon_{nk}$$

$$\frac{d\hat{s}}{dt} = \frac{\partial \hat{s}}{\partial t} + \dot{\mathbf{r}} \cdot \frac{\partial \hat{s}}{\partial \mathbf{r}}$$

$$\beta = \sum_{nk} \alpha_{nk} \frac{\mathbf{v}_{nk}}{\mathbf{v}_s}$$





Newton's apple tree
England



Descendant of
Newton's apple tree
NIST, Maryland

HAPPY BIRTHDAY!

**Thanks for the chaotic
upbringing**



END

Damped expectations for the second generation

